



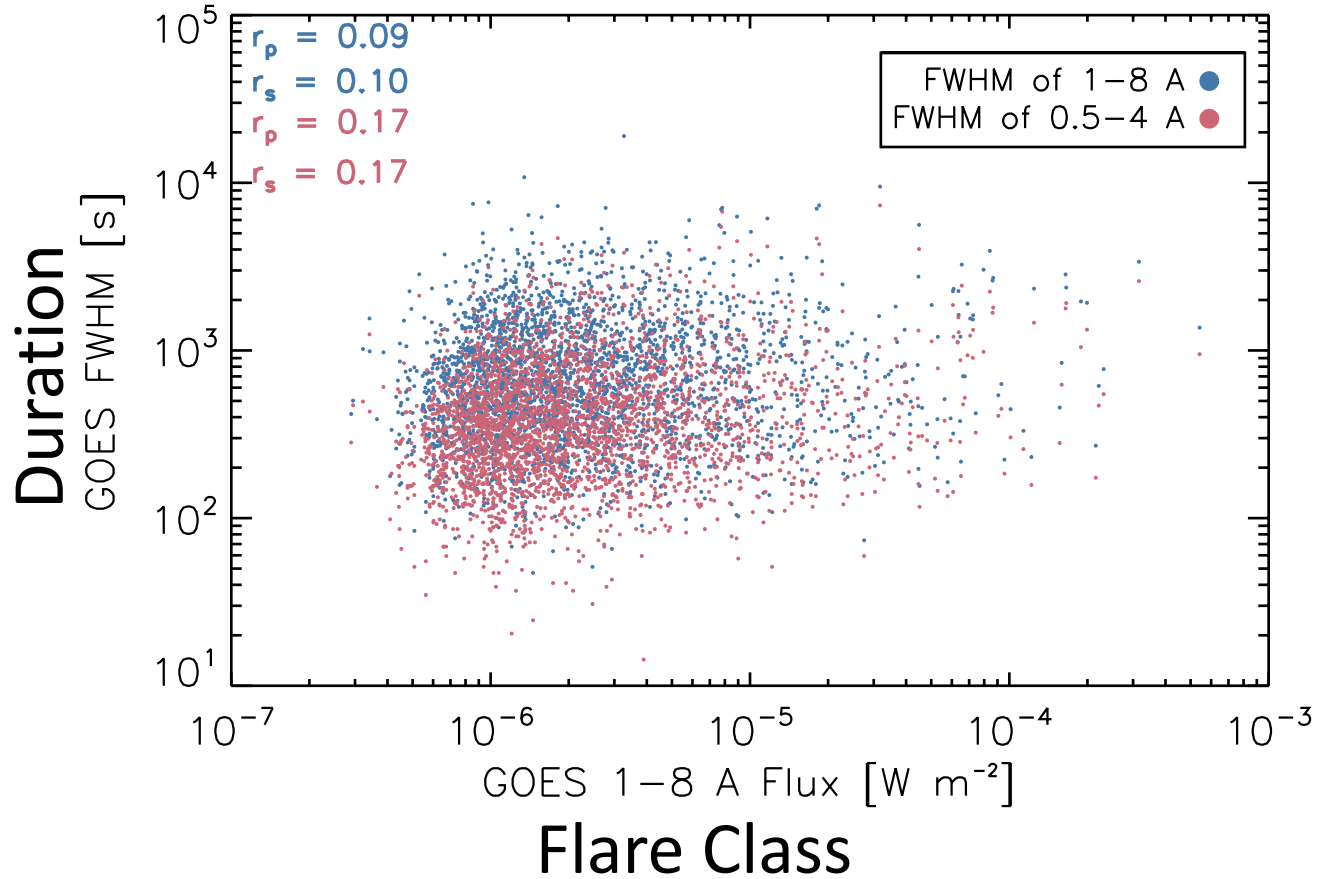
# Using Intensity/Duration Correlations in Solar and Stellar Flares to Improve Models of Irradiance Variability

Jeffrey W. Reep<sup>1</sup>

Vladimir Airapetian<sup>2,3</sup>, Sherry Chhabra<sup>4</sup>, Harry P. Warren<sup>1</sup>

<sup>1</sup>Space Science Division, US Naval Research Laboratory, Washington, DC, <sup>2</sup>American University, Washington, DC, <sup>3</sup>NASA Goddard Space Flight Center, Greenbelt, MD <sup>4</sup>George Mason University, Fairfax, VA

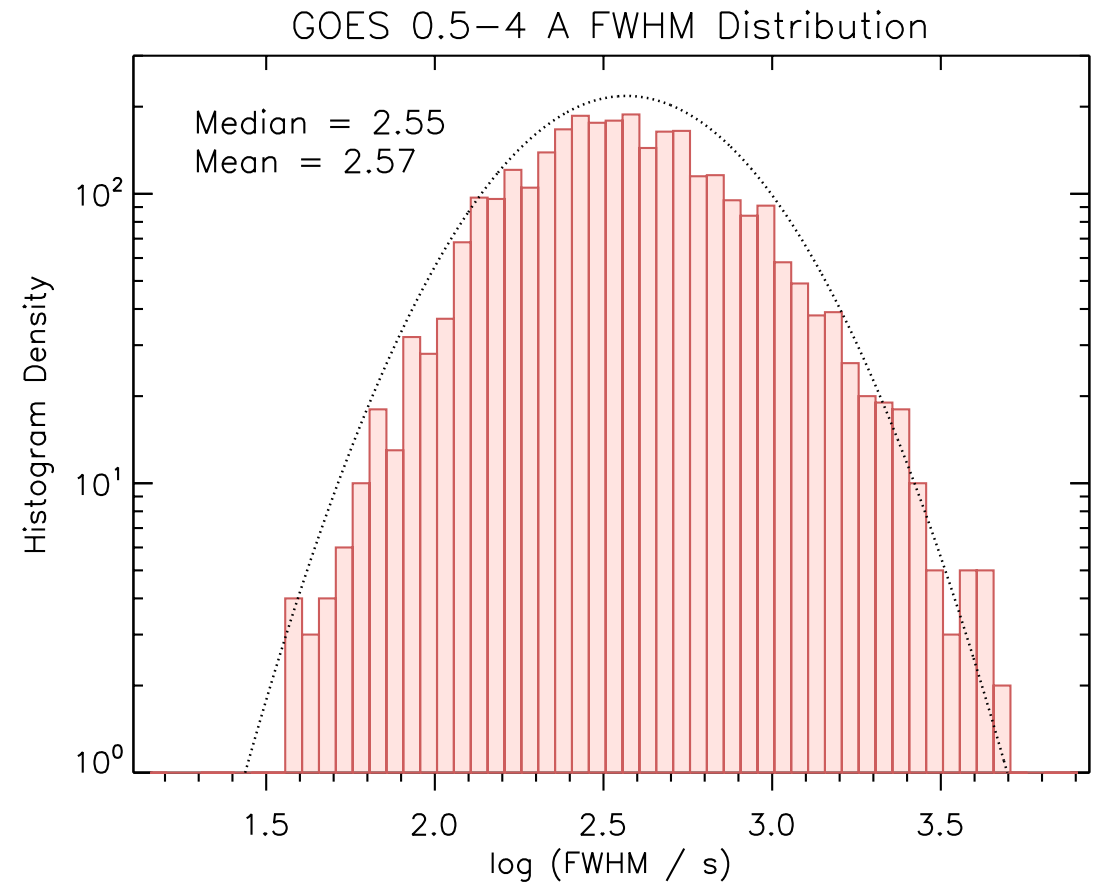
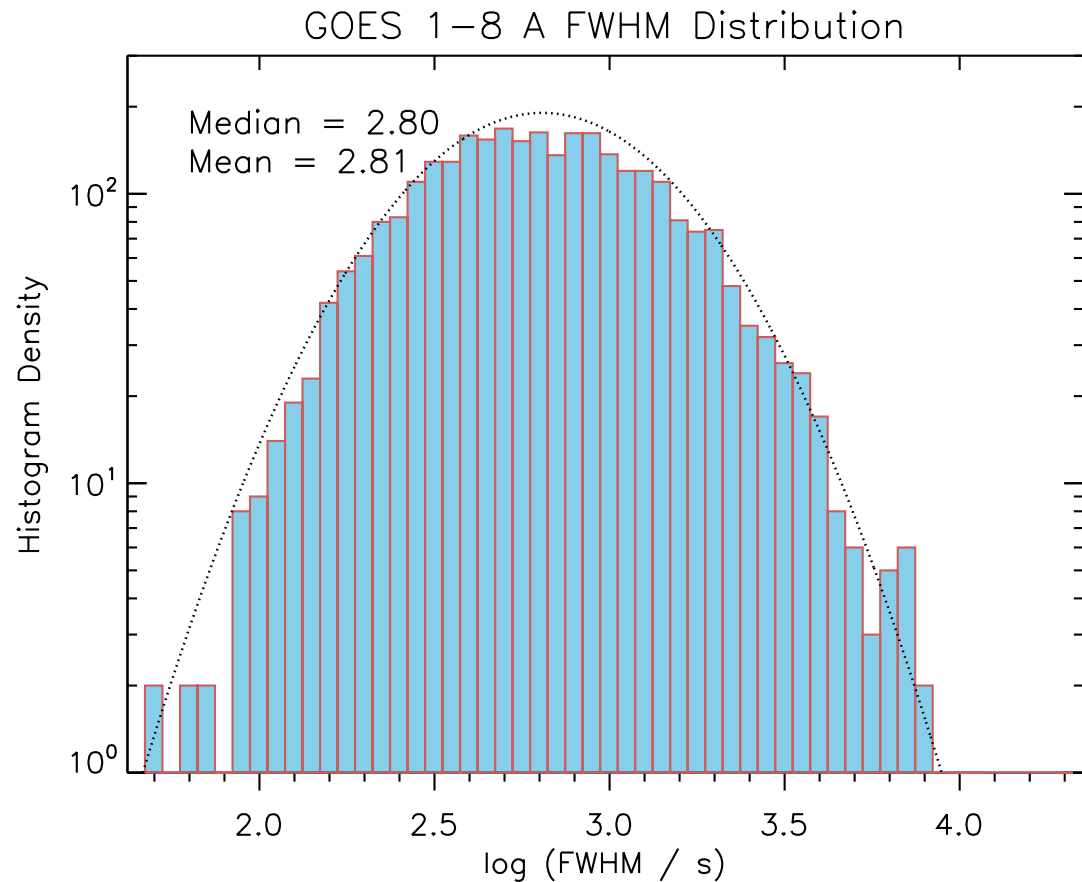




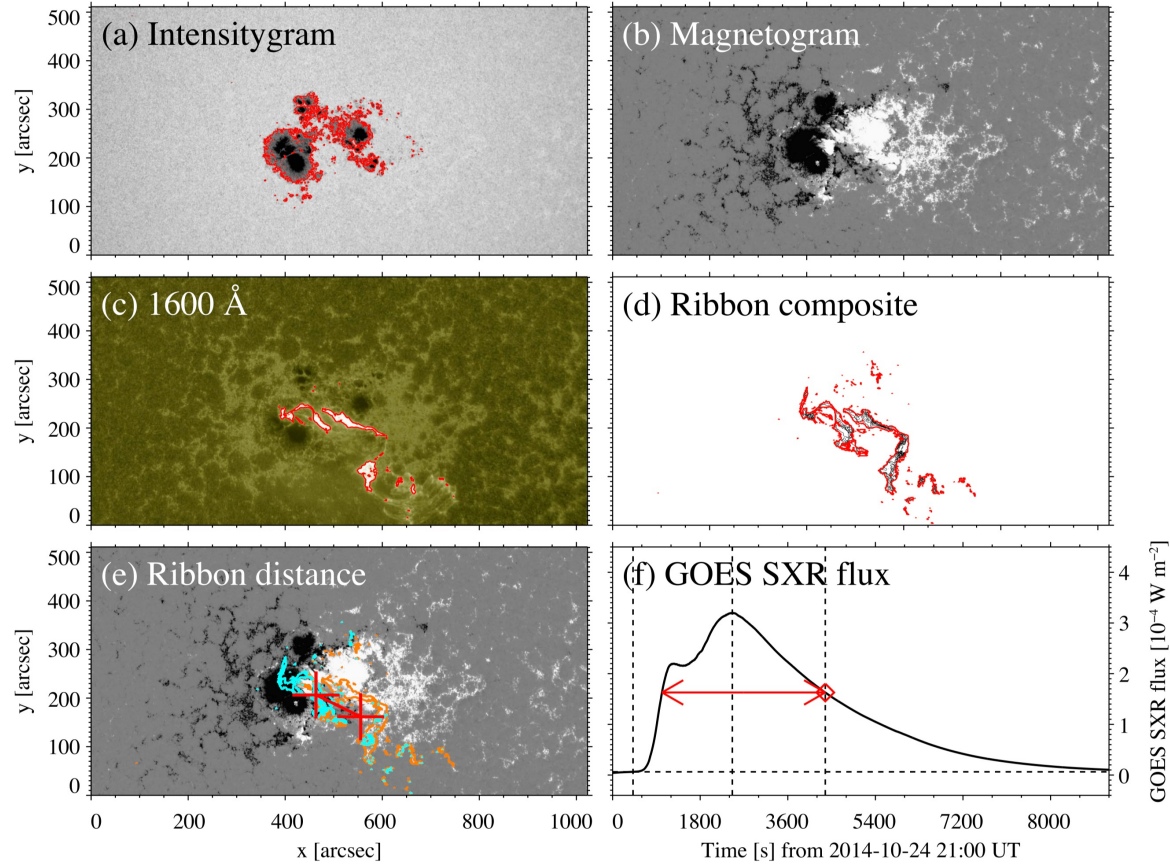
**Not correlated!**

Duration also not correlated with physical volume, temperature, density, total energy release, or magnetic field strength

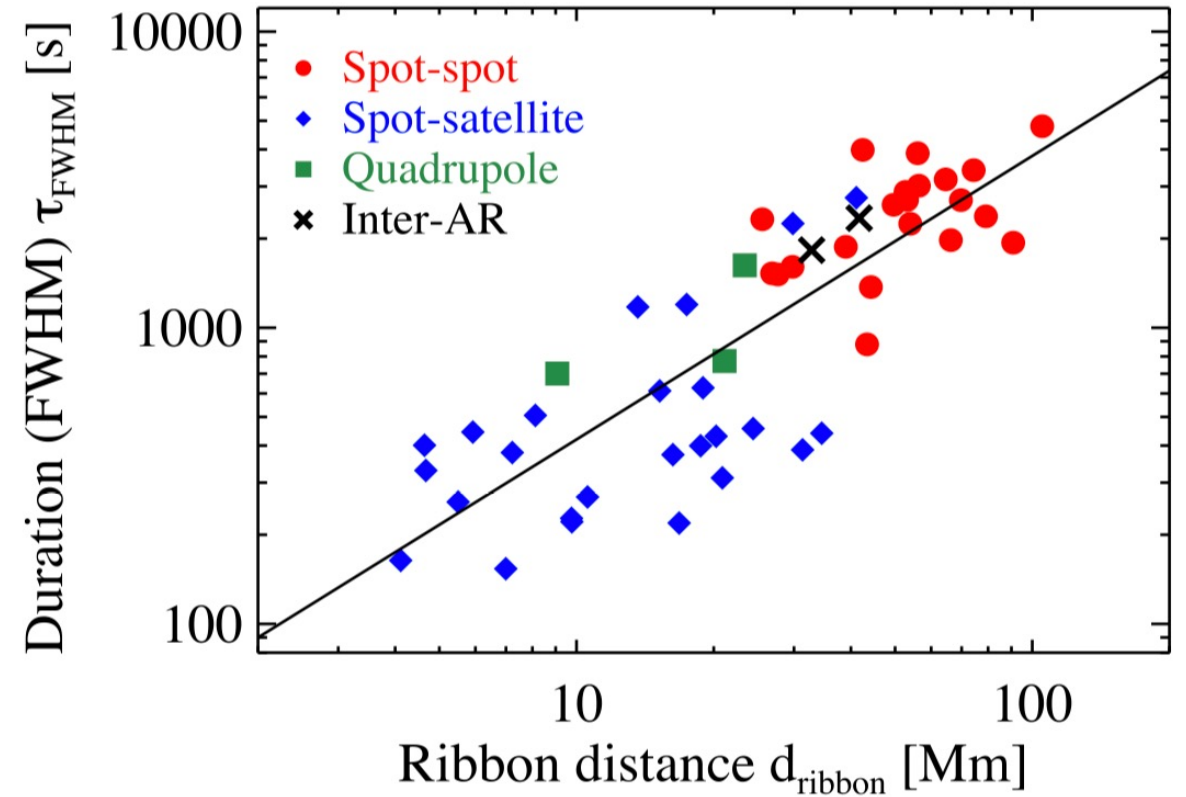
## Log-normal Distributions, dependent on wavelength



2014-10-24 X3.1-class flare NOAA AR 12192

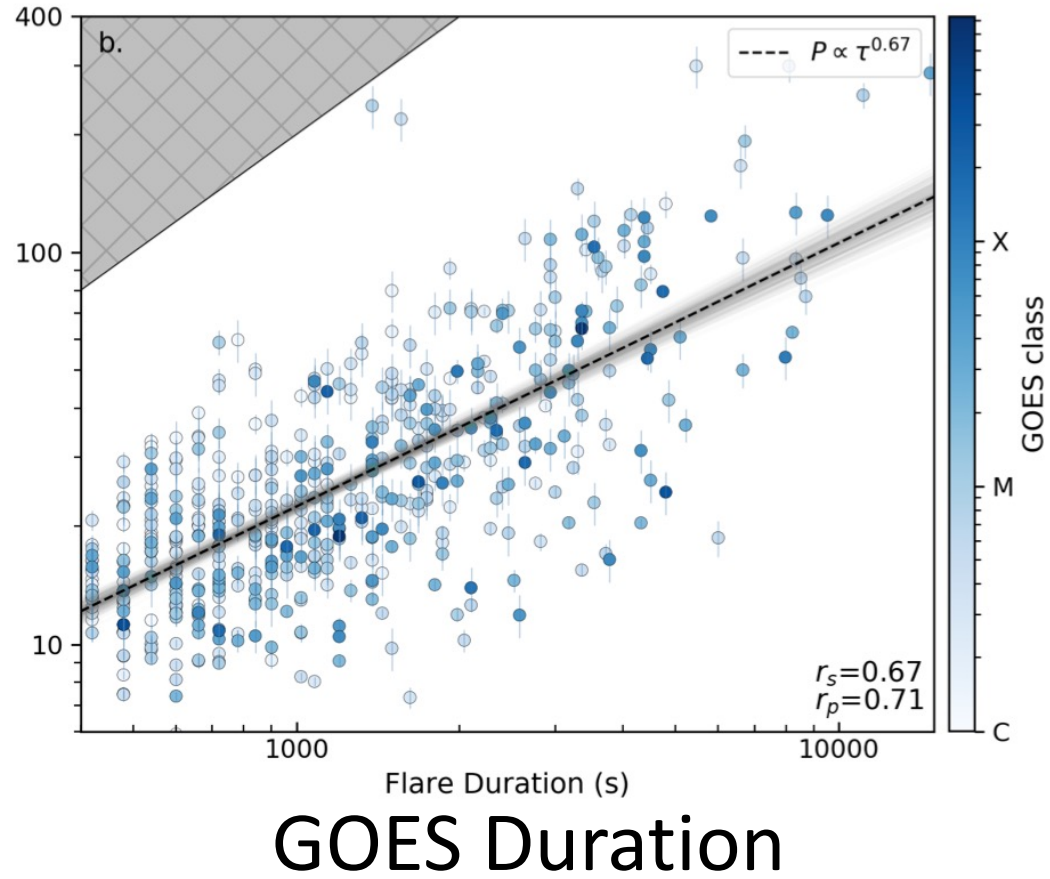
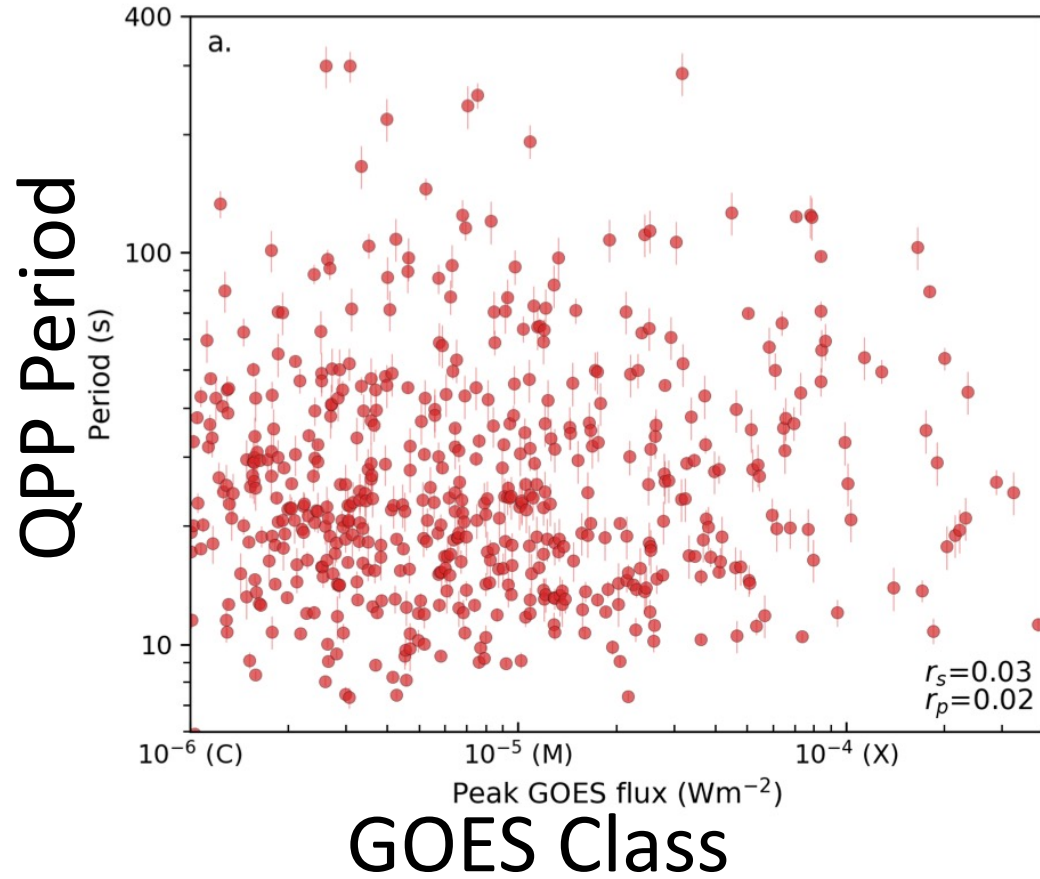


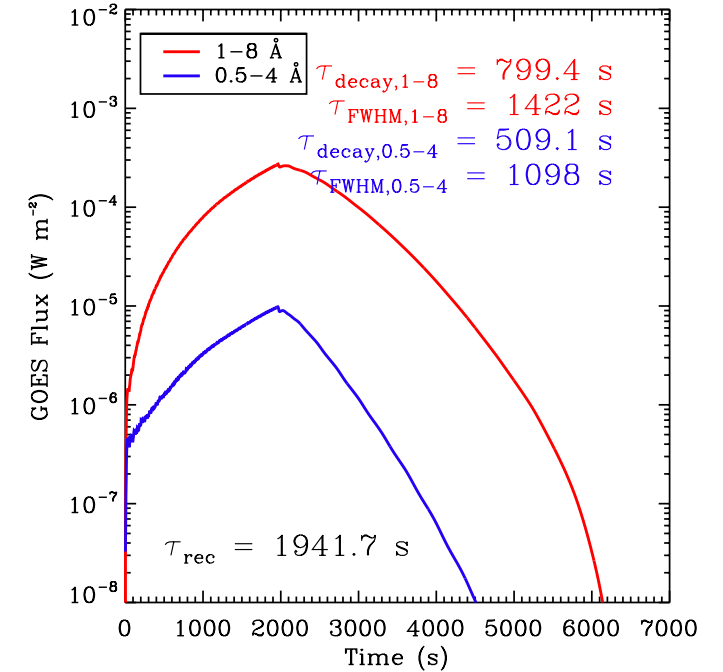
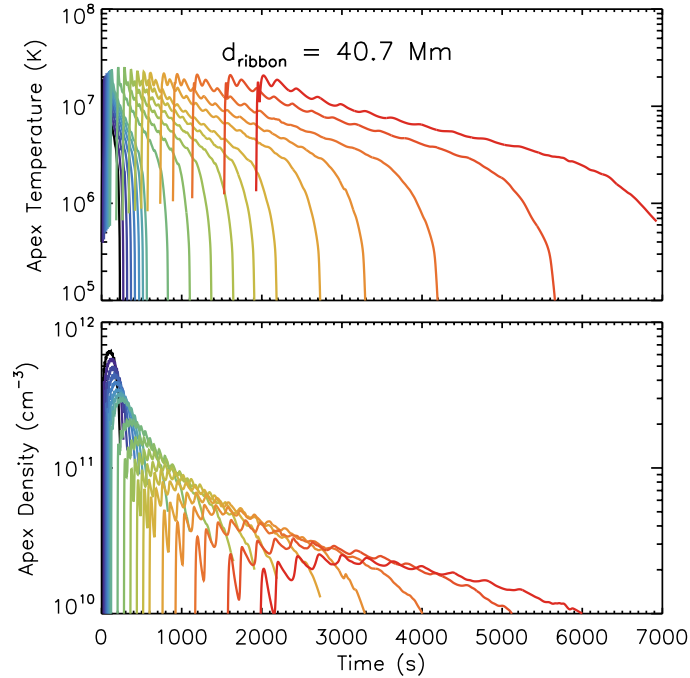
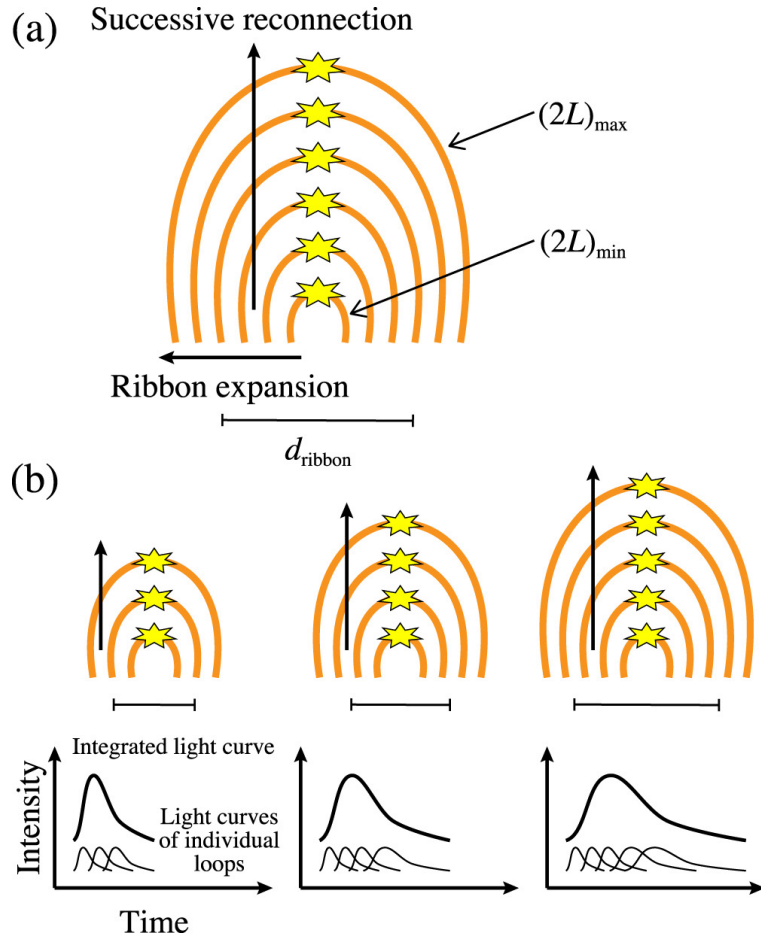
## Ribbon separation correlated with flare duration



# QPP period related to SXR duration, but not class

Reep et al.

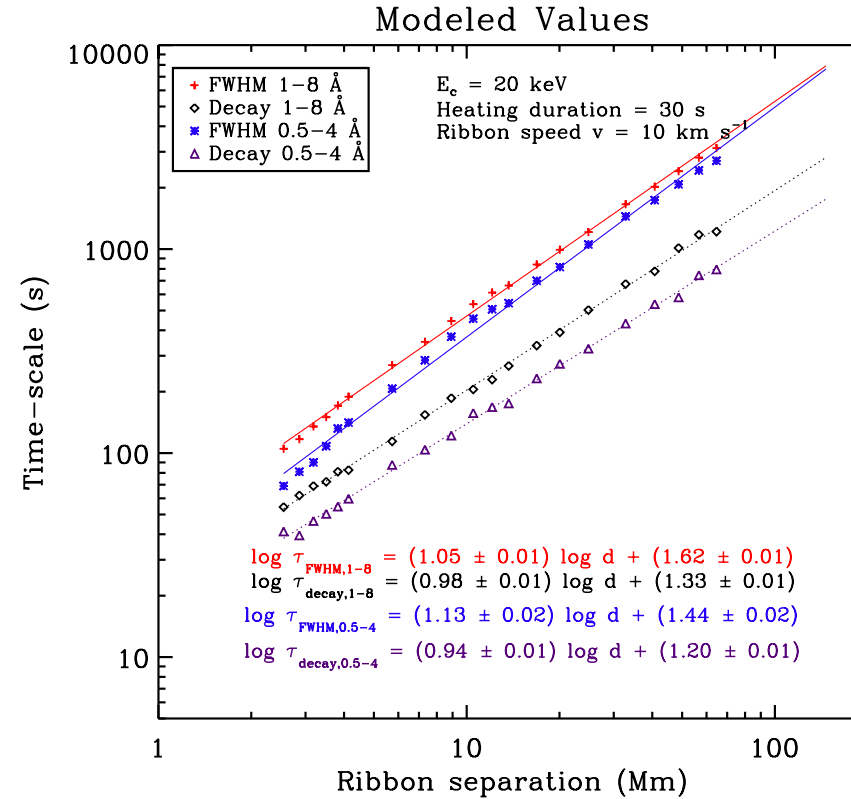
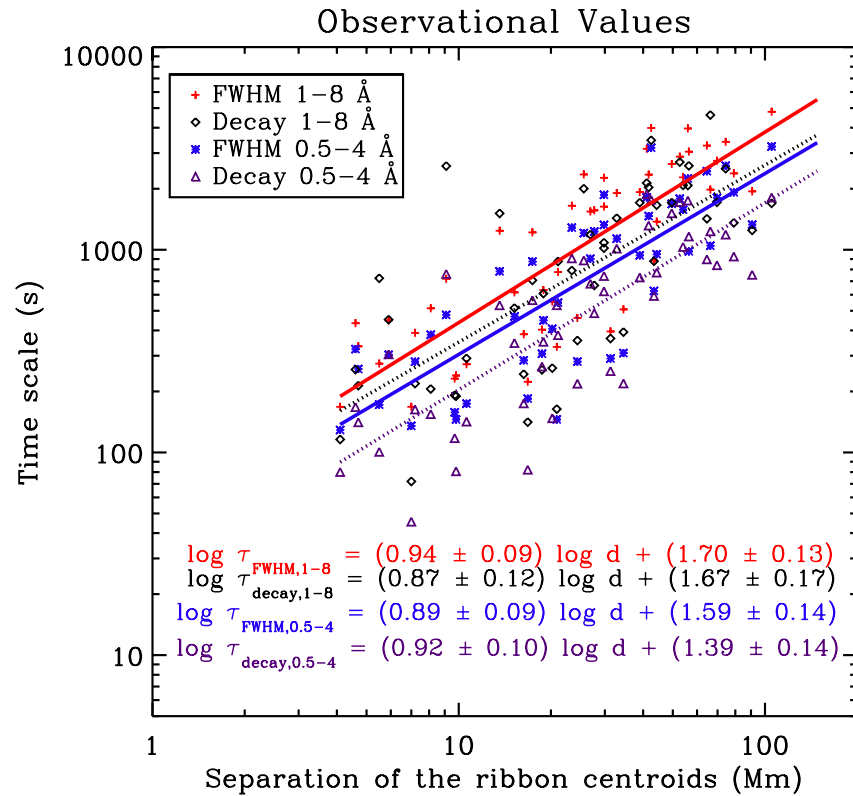




Construct a multithreaded flare model to test the relation

# $\tau$ - d Relation Reproduced

Reep et al.



Reconnection duration can explain the relation between  
SXR duration and ribbon separation

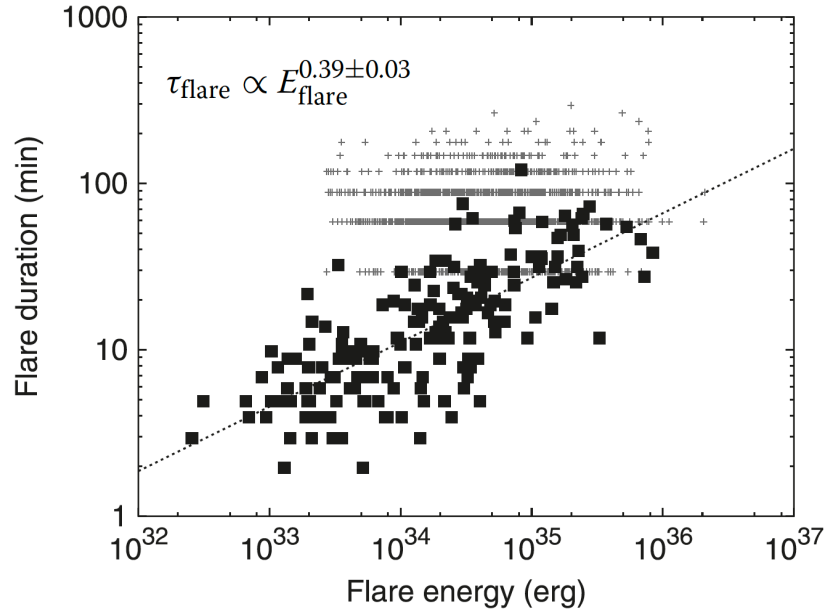
Likely explanation is that the duration of a flare is simply a measure of the time to convert the magnetic flux to thermal energy

- More flux, longer time reconnecting to longer and longer loops
- QPP period tied to loop length ( $\sim L/v_A$ )

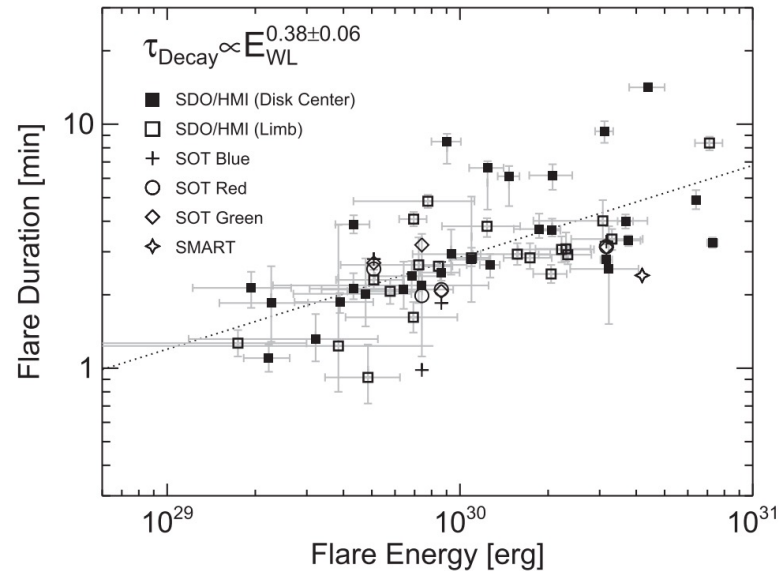
Are we done?



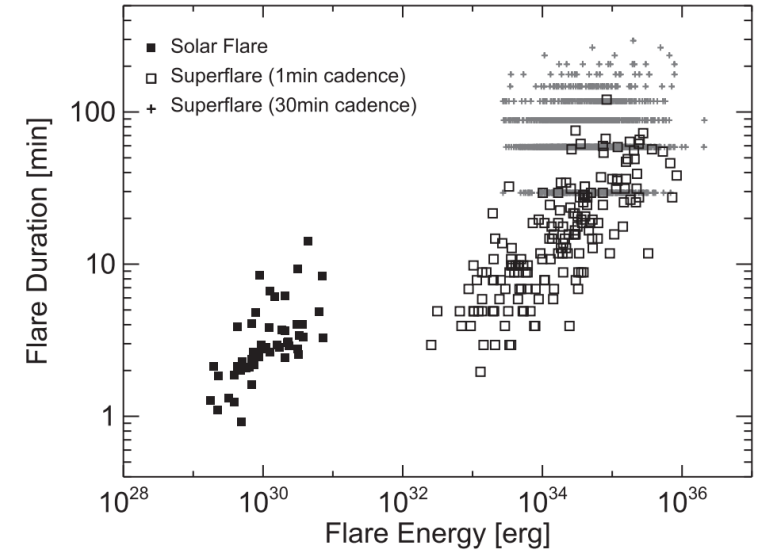
Stellar flares observed by Kepler



Solar flares observed by HMI and SOT

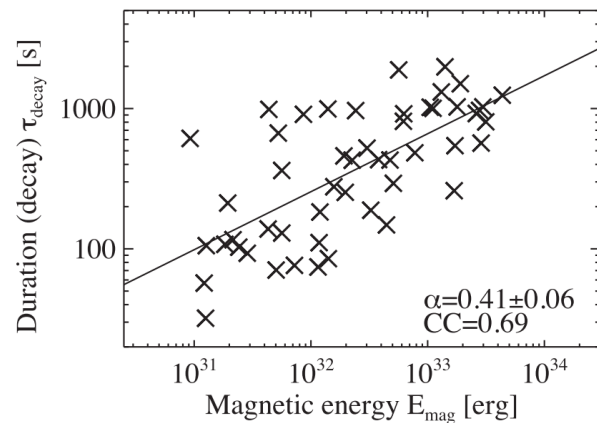
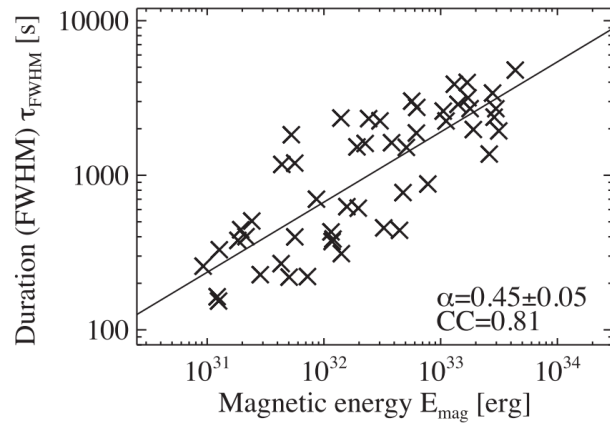


Both



- Both solar and stellar flares show the same relation between white light duration and flare energy
- Stellar flares are longer duration, possibly because of higher field strengths (longer magnetic dissipation)

Time versus magnetic energy



SXR duration *is* correlated with magnetic energy of the event

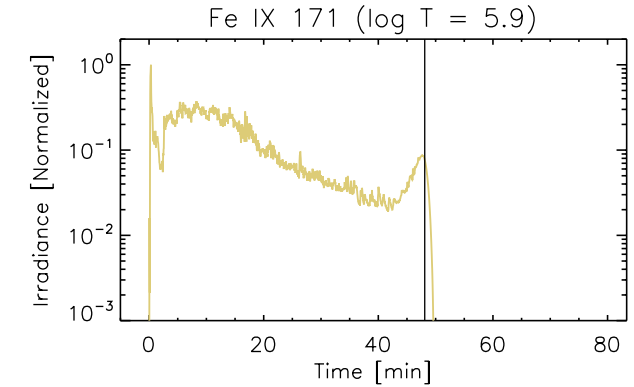
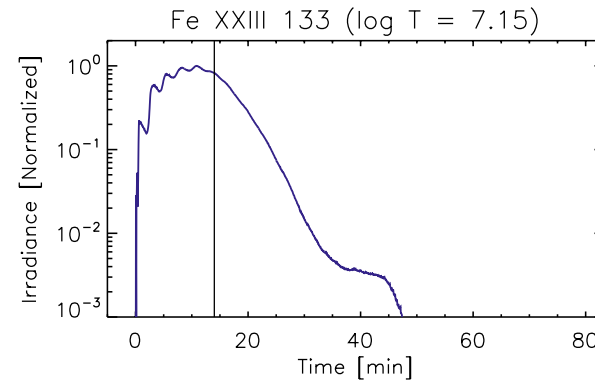
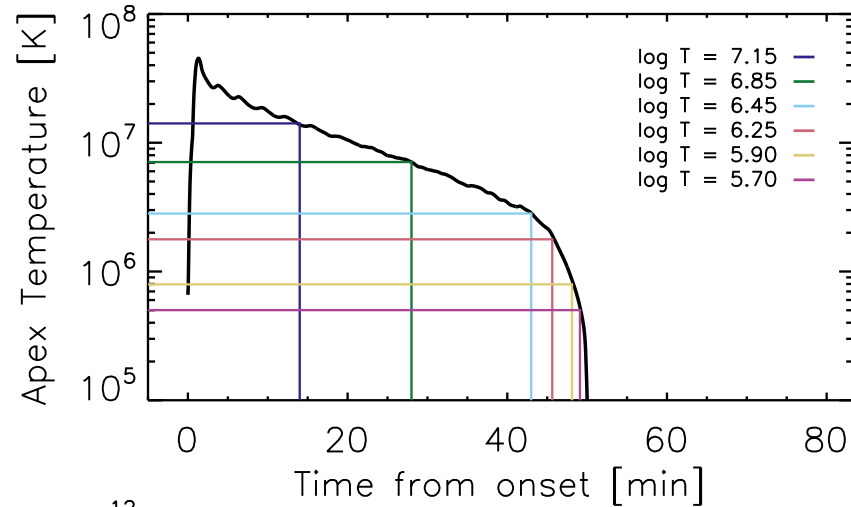
→ There is some sort of disconnect between magnetic energy and emergent X-ray intensity

# Irradiance time series at different temperatures

## 1. Uniform, laminar loop

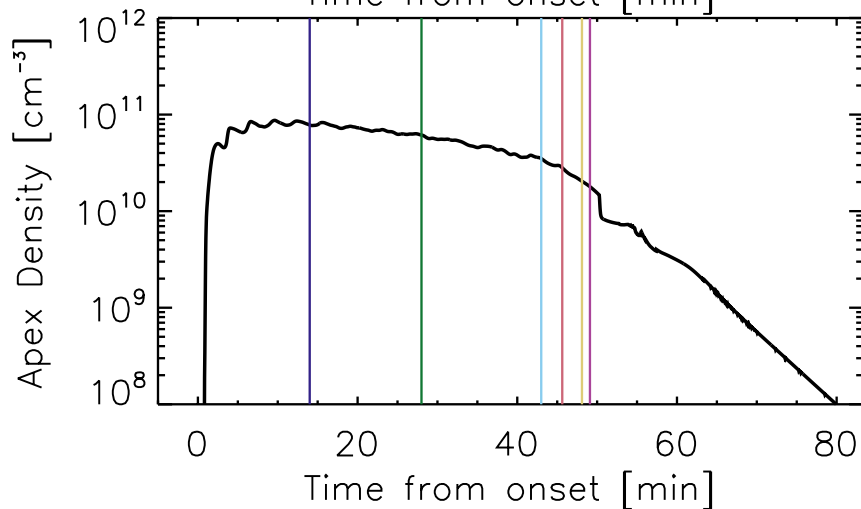
Reep et al.

$$\tau_{cool} \propto \frac{L^{5/6}}{(nT)^{1/6}}$$



In a "typical" flaring loop, emission forming at hot and cool temperatures have a distinct time evolution

→ The duration of emission in a given wavelength depends on where it forms in the atmosphere

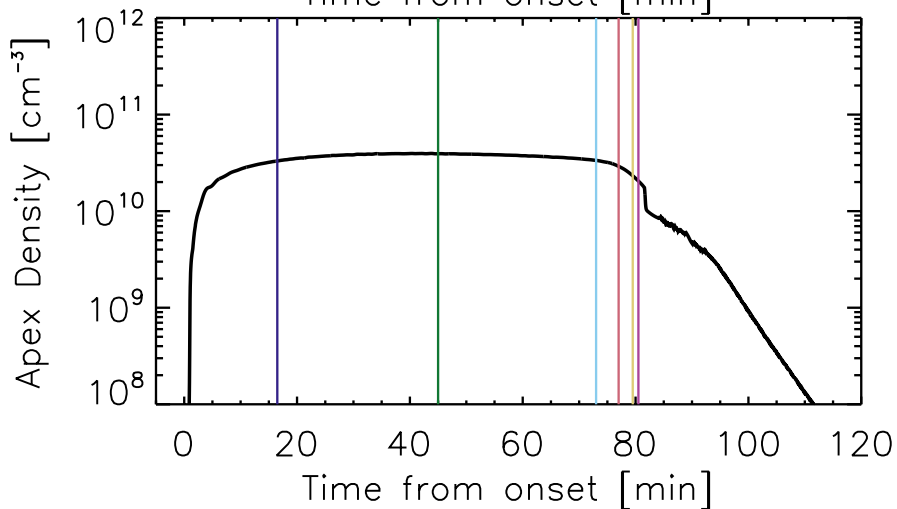
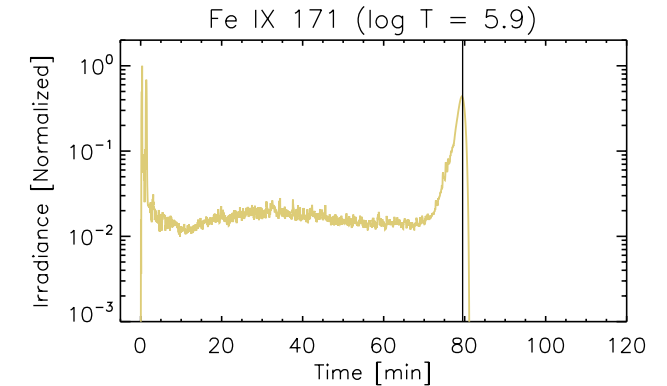
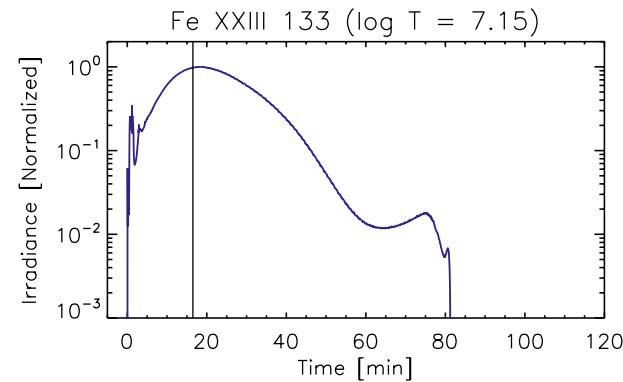
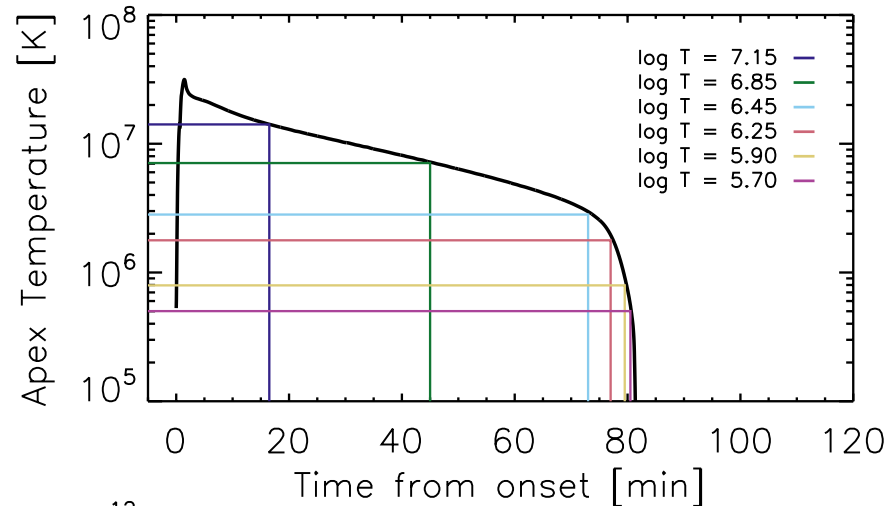


# Irradiance time series at different temperatures

## 2. Expanding, laminar loop

Reep et al.

$$\tau_{cool} \propto ?$$



With area expansion, the cooling time is lengthened, and draining is suppressed which causes light curves to track ionization fractions

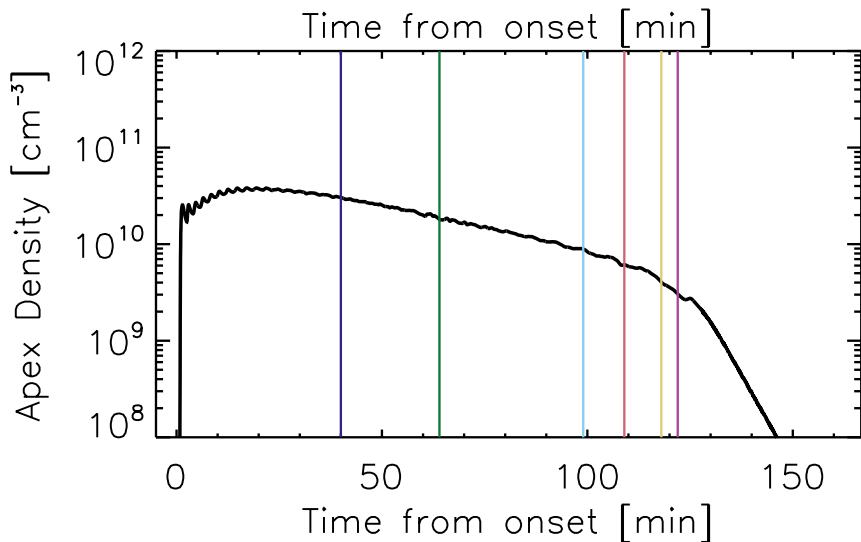
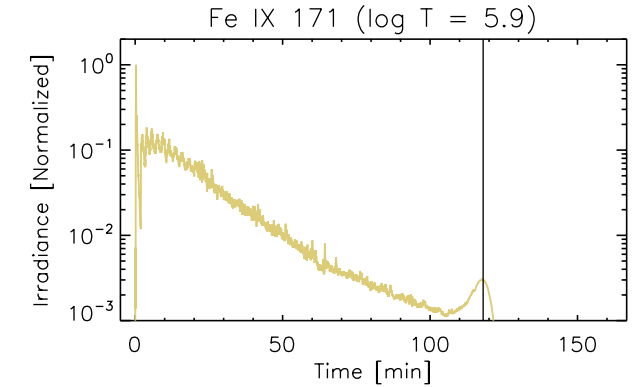
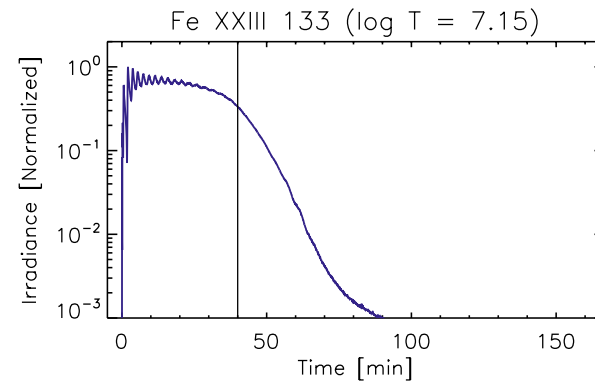
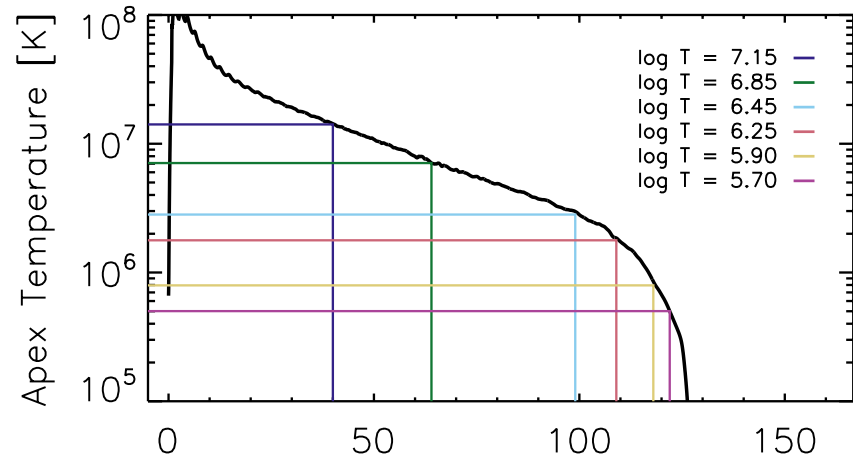
→ The duration of emission depends on the change in magnetic topology (equivalently, field strength)

# Irradiance time series at different temperatures

## 3. Uniform, turbulent loop

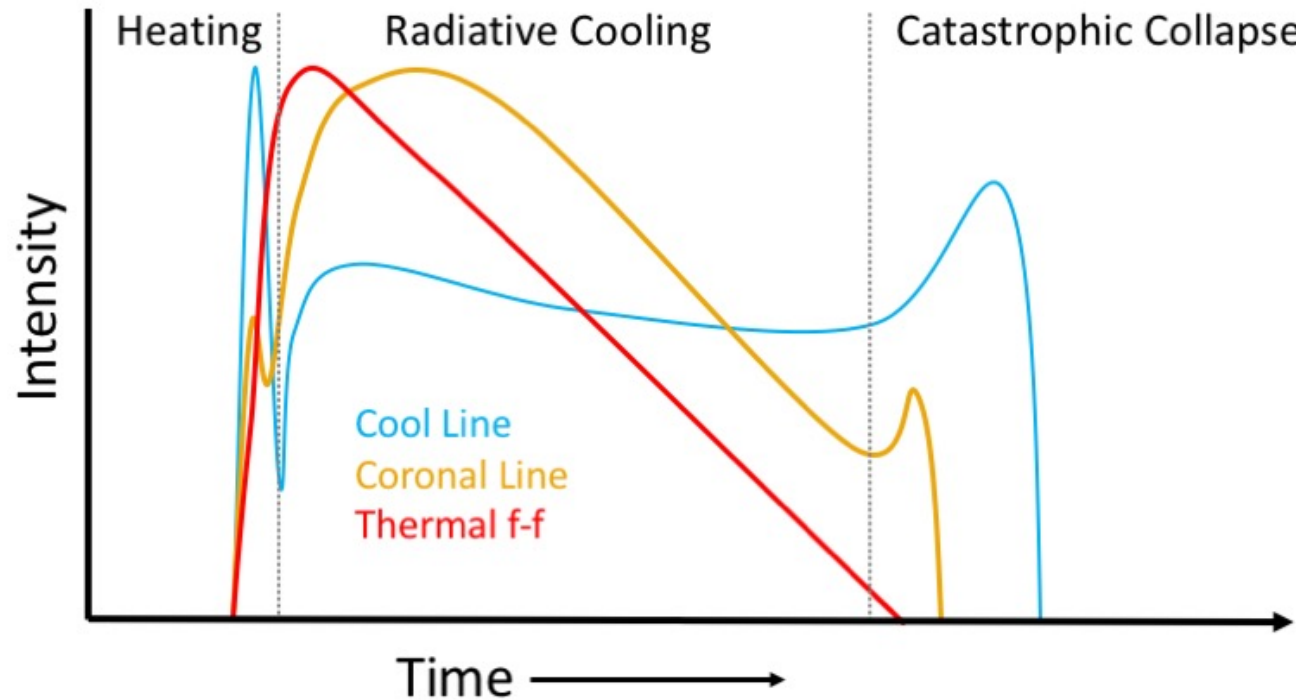
Reep et al.

$$\tau_{cool} \propto \frac{T^{2/3} L^{5/6}}{\lambda_T n^{7/12}}$$



With turbulence, the cooling time is significantly lengthened, and the loop steadily drains. Light curves slowly decrease with time.

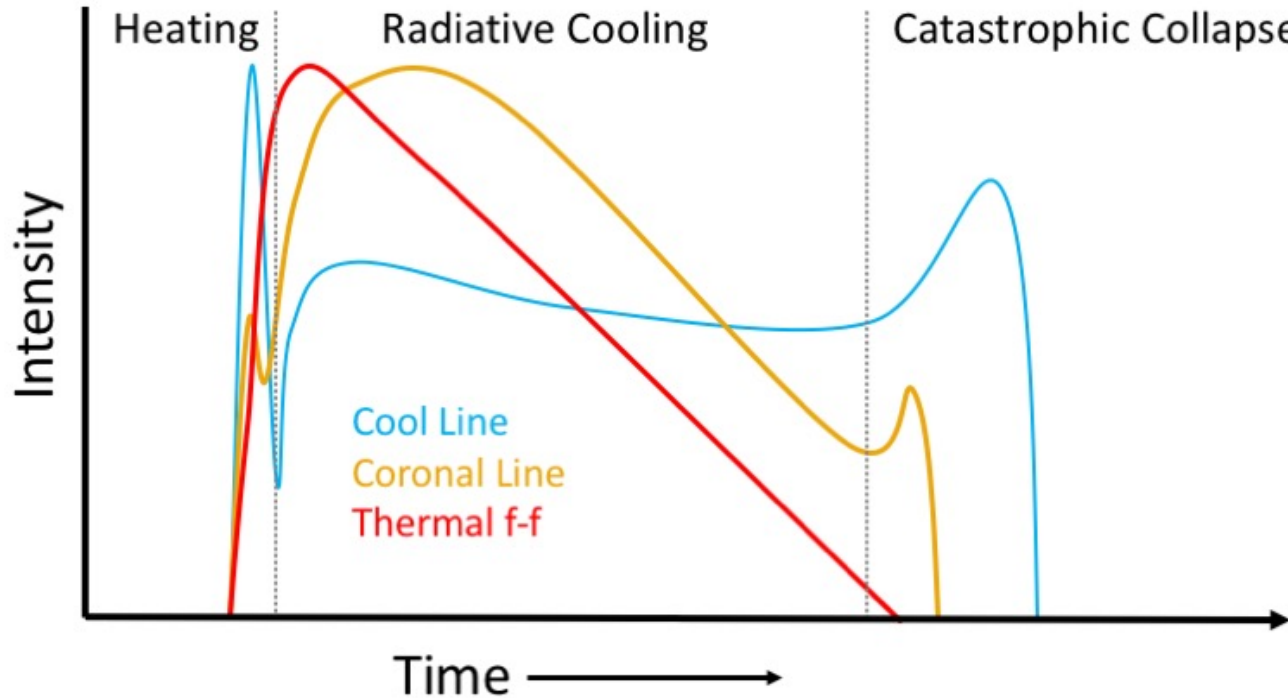
→ The duration of emission depends on how turbulent the plasma is



Three distinct phases:

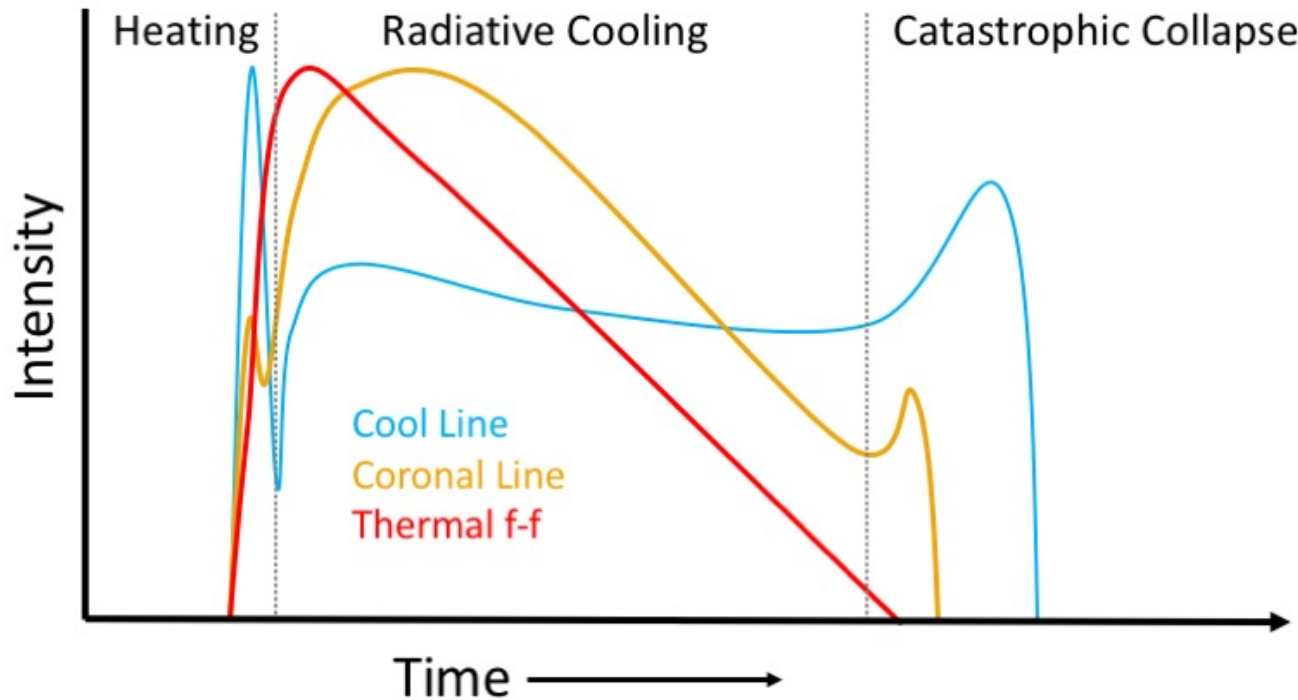
1. Heating phase, where all emission rises in step with the heating. Cool emission spikes strongly, as the transition region is lowered in height to higher densities.

(We saw in Harry's talk yesterday that Ly- $\alpha$  spiked simultaneously with HXR's, for example)



Three distinct phases:

2. Radiative cooling phase, where the density of the corona peaks due to evaporation, while the temperature starts to fall. Cooler emission, forming in the TR or chromosphere, is mostly steady during this time. Hotter emission, forming in the corona, decays as the temperature falls.



Three distinct phases:

3. Catastrophic collapse, when the coronal temperatures plummets to chromospheric values, and the density quickly drains. Individual lines spike in intensity when the coronal portion of the loop reaches their peak formation temperature, but then fade away almost instantly.



There is an obvious distinction between the behavior of cool and hot emission, and this stems from the so-called Neupert effect, originally a relation found between thermal SXR emission and non-thermal microwave or hard X-ray emission:

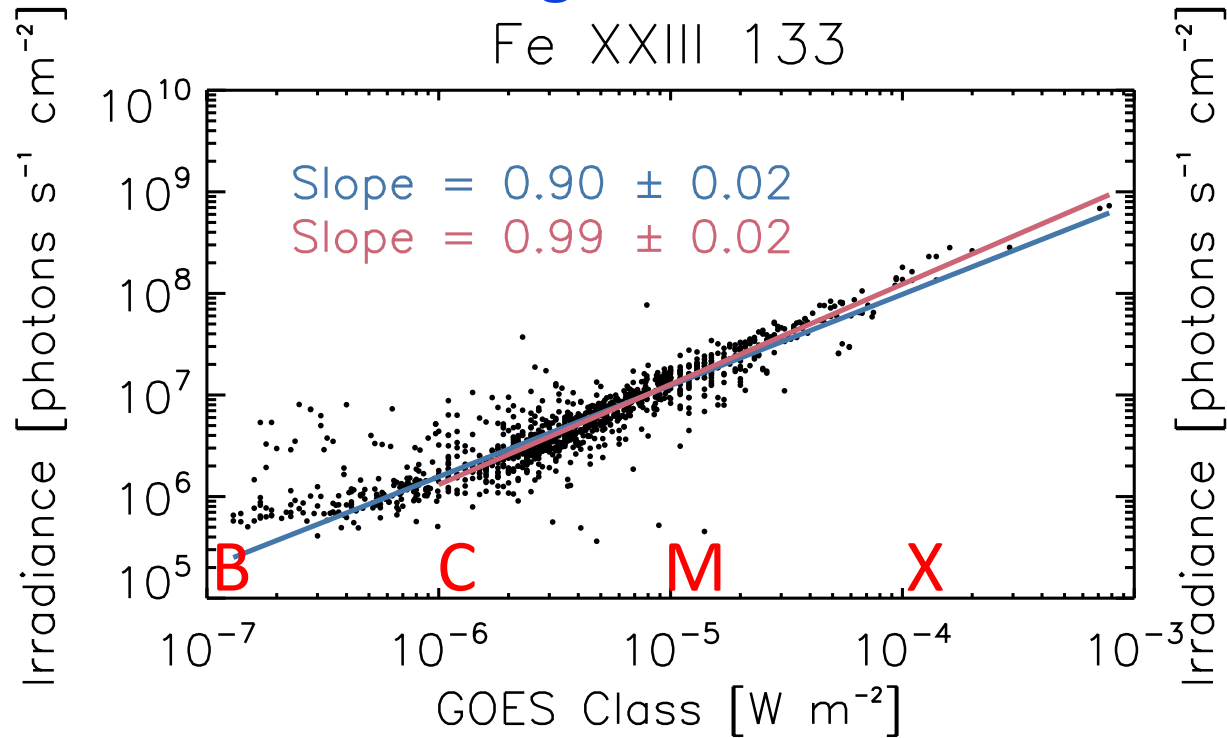
$$I_{\text{SXR}} \propto \int I_{\text{HXR}} dt$$

$$I_{\text{HXR}} \propto \frac{d}{dt} (I_{\text{SXR}})$$

The relation is more general. Emission forming in the low atmosphere (chromosphere, TR) responds directly to heating as the TR is pushed to lower heights and higher densities. Emission forming in the corona increases with chromospheric evaporation, that is, indirectly to heating.

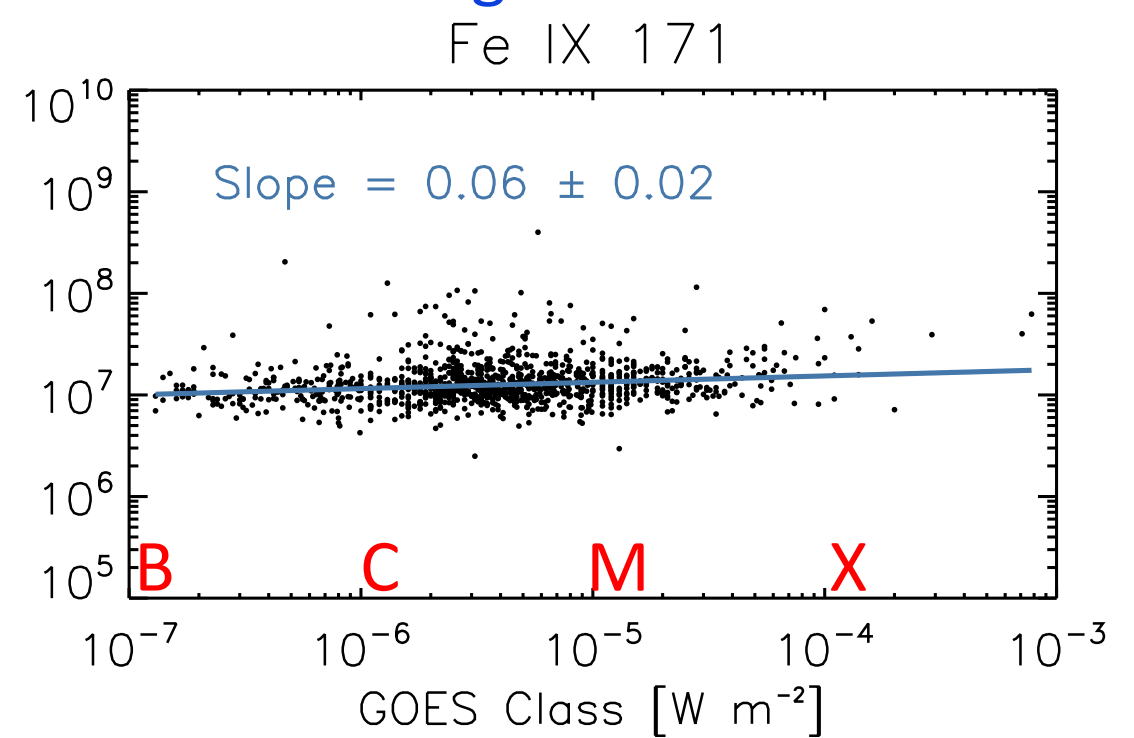
$\log T = 7.15$

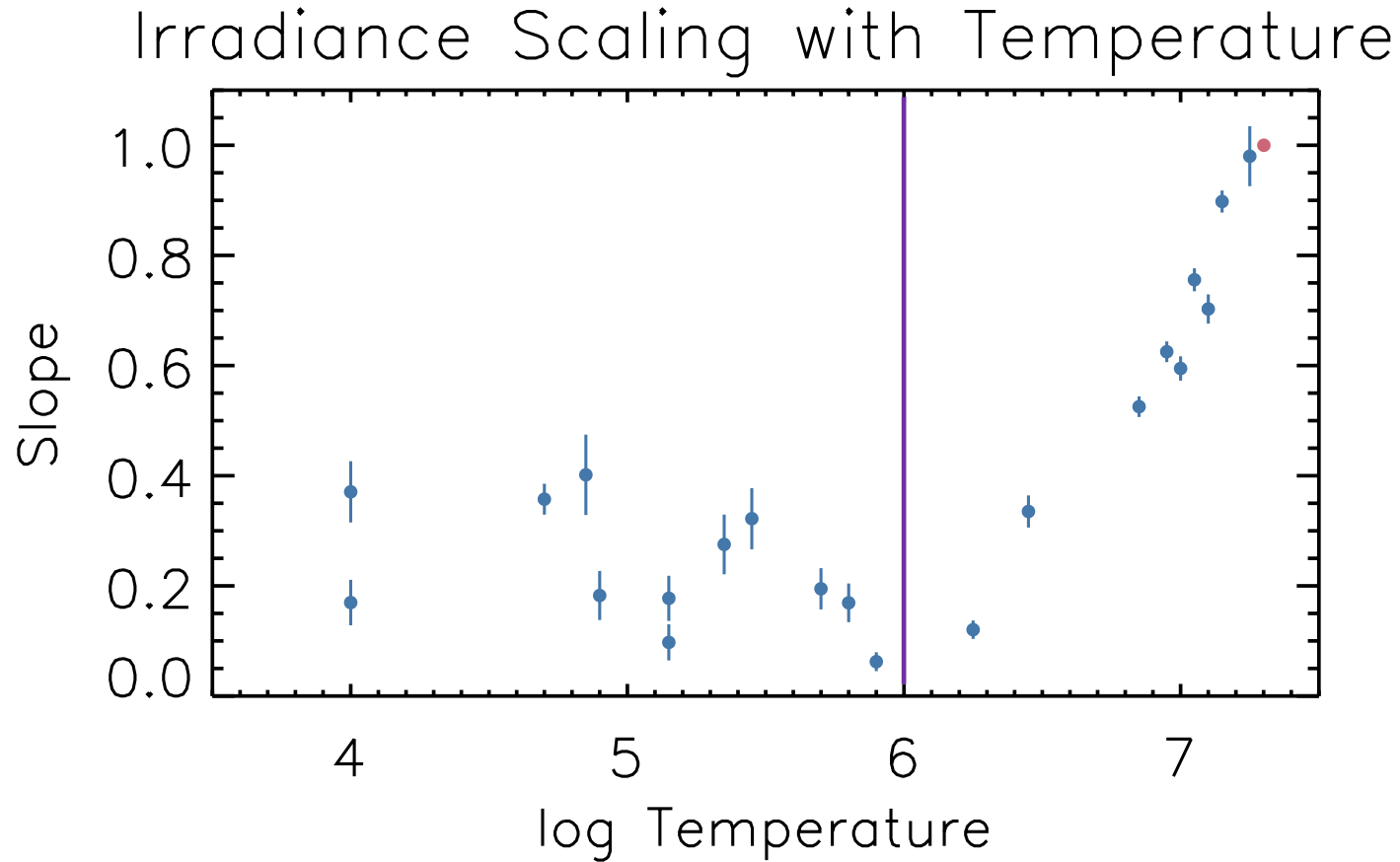
Fe XXIII 133



$\log T = 5.90$

Fe IX 171





Hot lines scale  $\sim$  linearly

Cooler lines scale sublinearly

Slope decreases with temperature until about  $\log T = 6$

The behavior of emission that responds directly to heating behaves differently to emission that responds indirectly (a general Neupert effect). SXR intensity is disconnected from the flare duration because it responds indirectly, brightening after the onset of chromospheric evaporation.

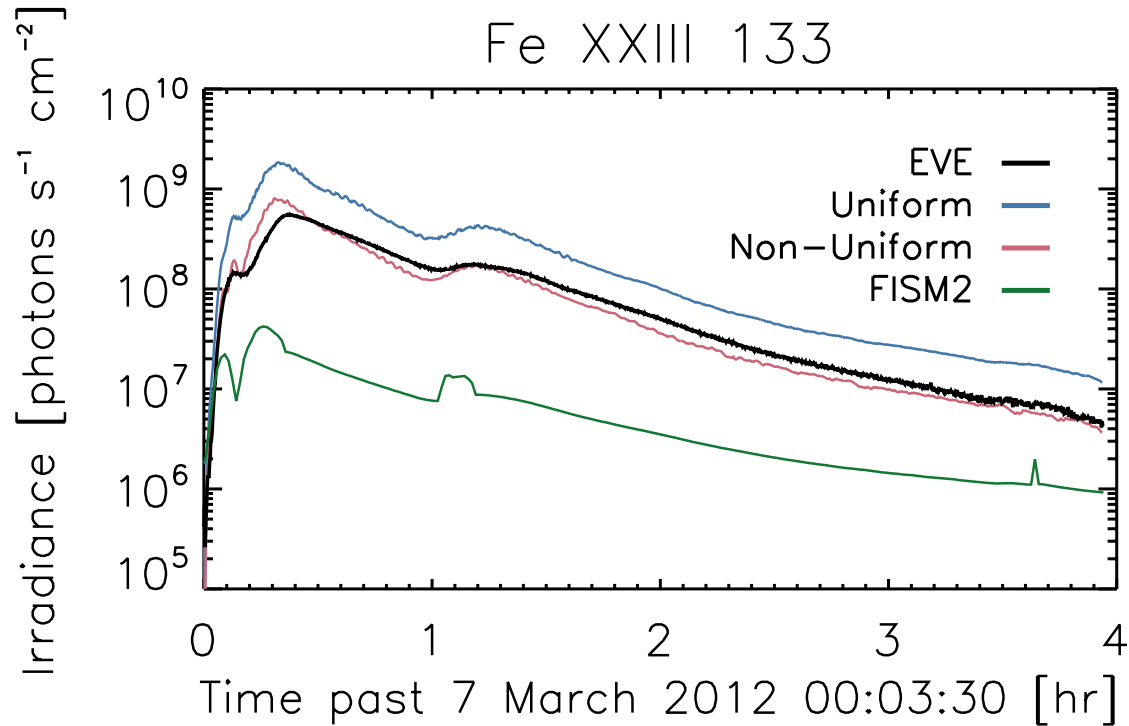
To improve models of irradiance, we need to:

- Understand the magnetic topology from the chromosphere through the corona
- Quantify the level of turbulence across the atmosphere

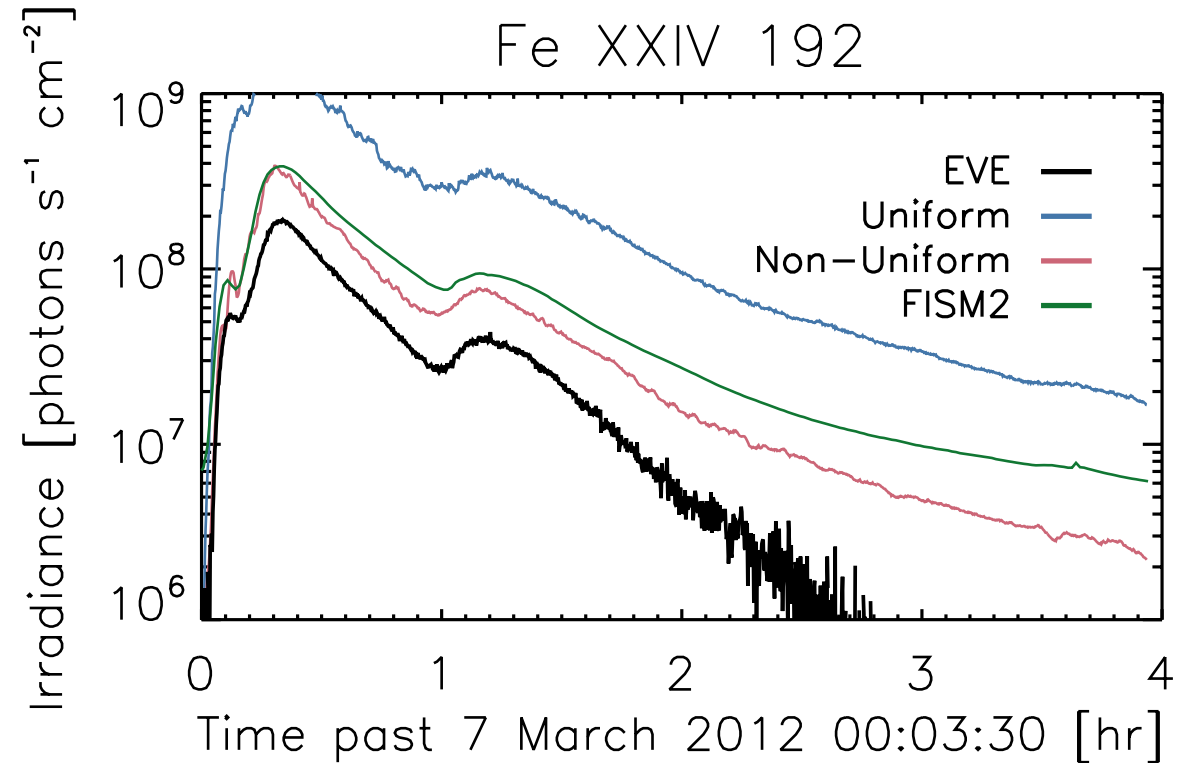
This work was supported by the Office of Naval Research 6.1 Support Program.

# Supplementary Slides

## 07-Mar-2012 UT 00:03 X7.8 Flare

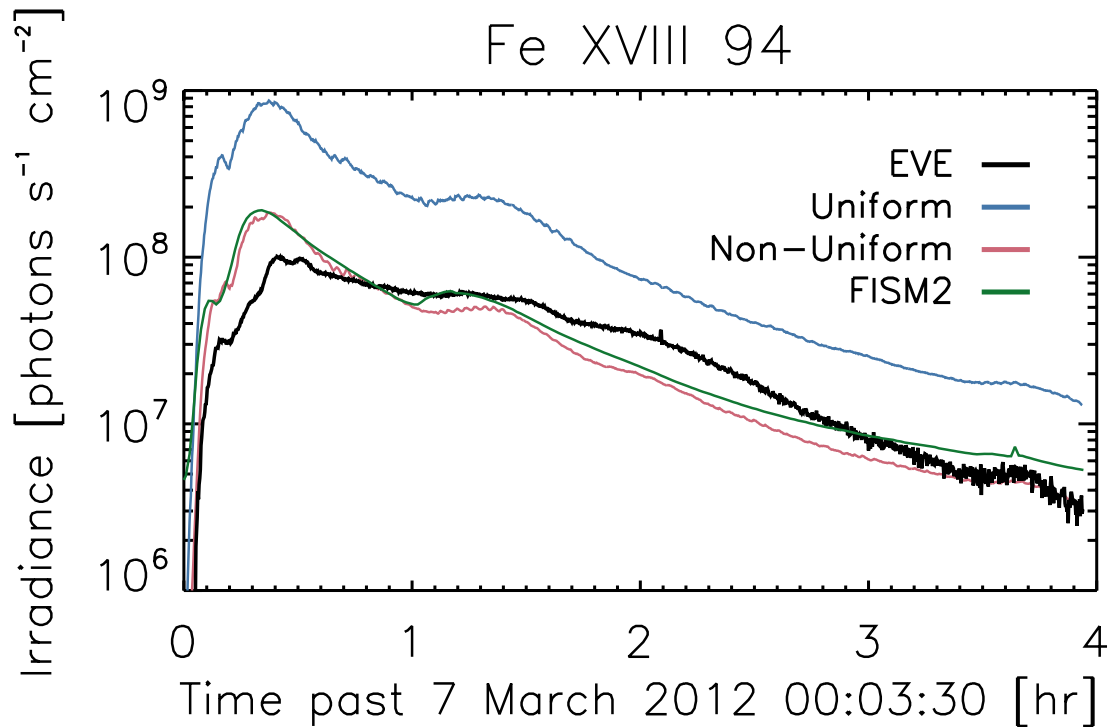


log T = 7.15

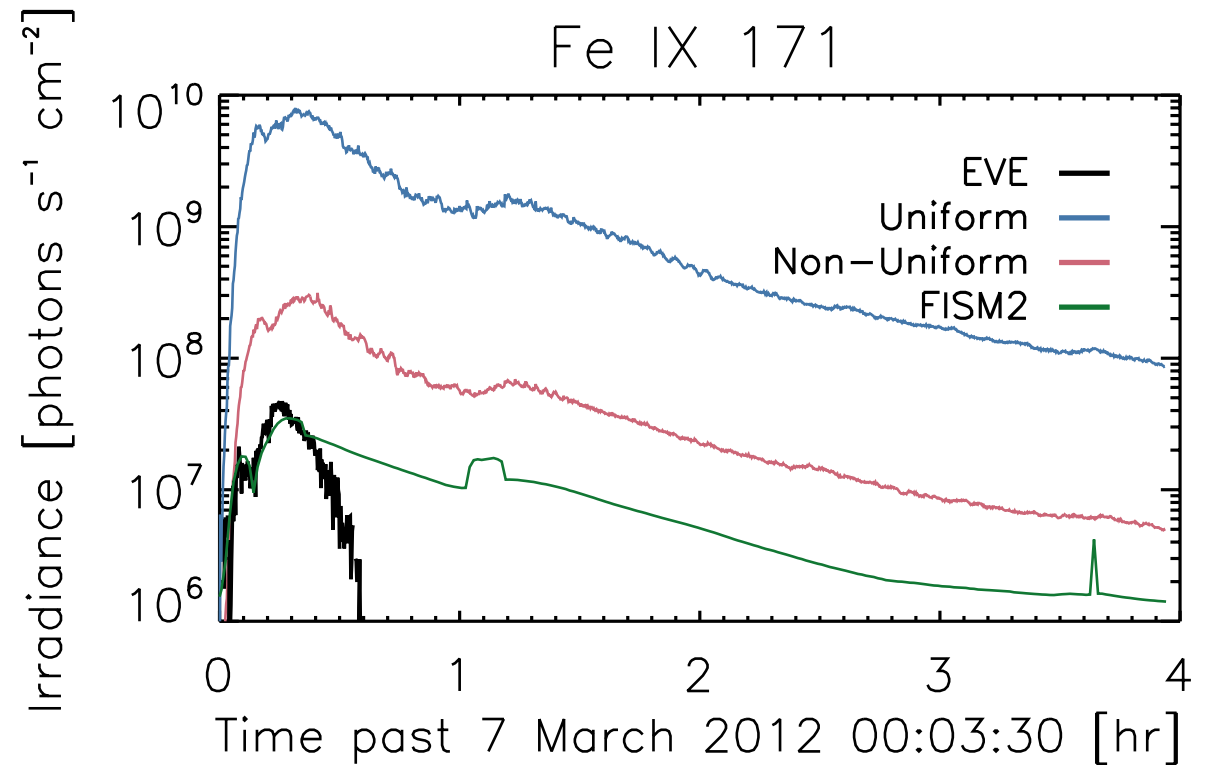


log T = 7.25

## 07-Mar-2012 UT 00:03 X7.8 Flare



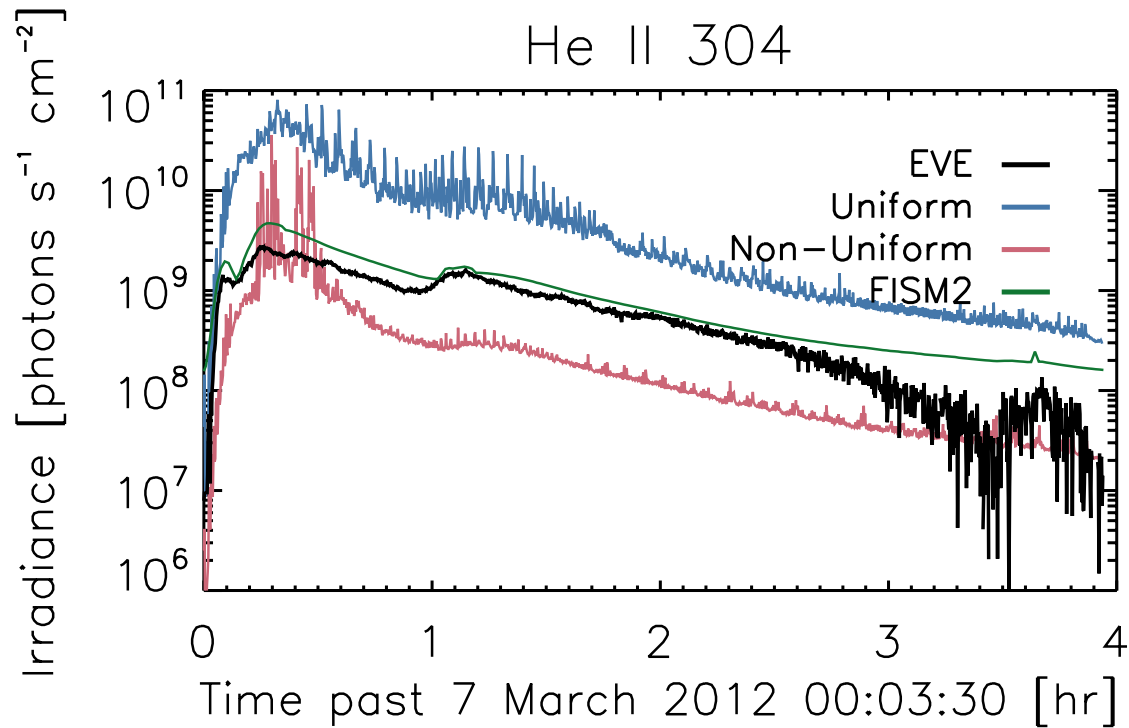
$\log T = 6.85$



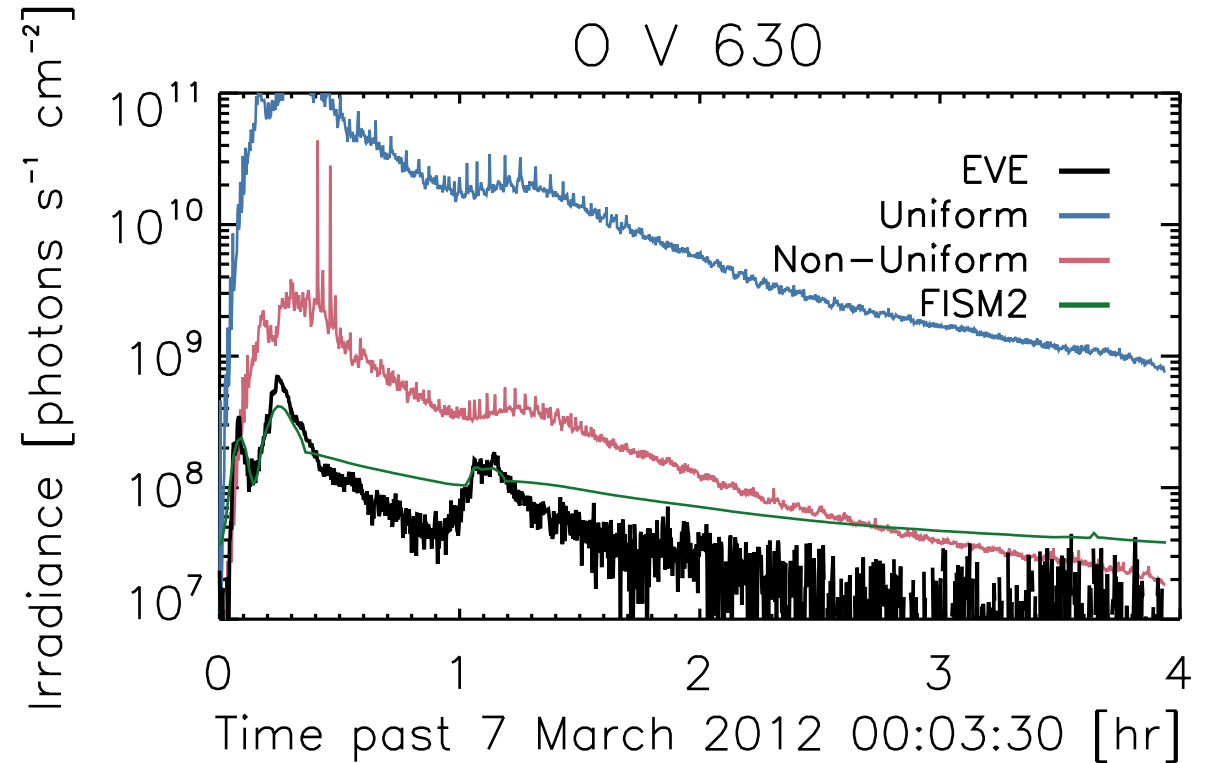
$\log T = 5.90$



## 07-Mar-2012 UT 00:03 X7.8 Flare



$\log T = 4.70$



$\log T = 5.35$

