

# Linear eigenmode analysis of the tearing instability in electron-scale current sheet with nongyrotropic electron pressure

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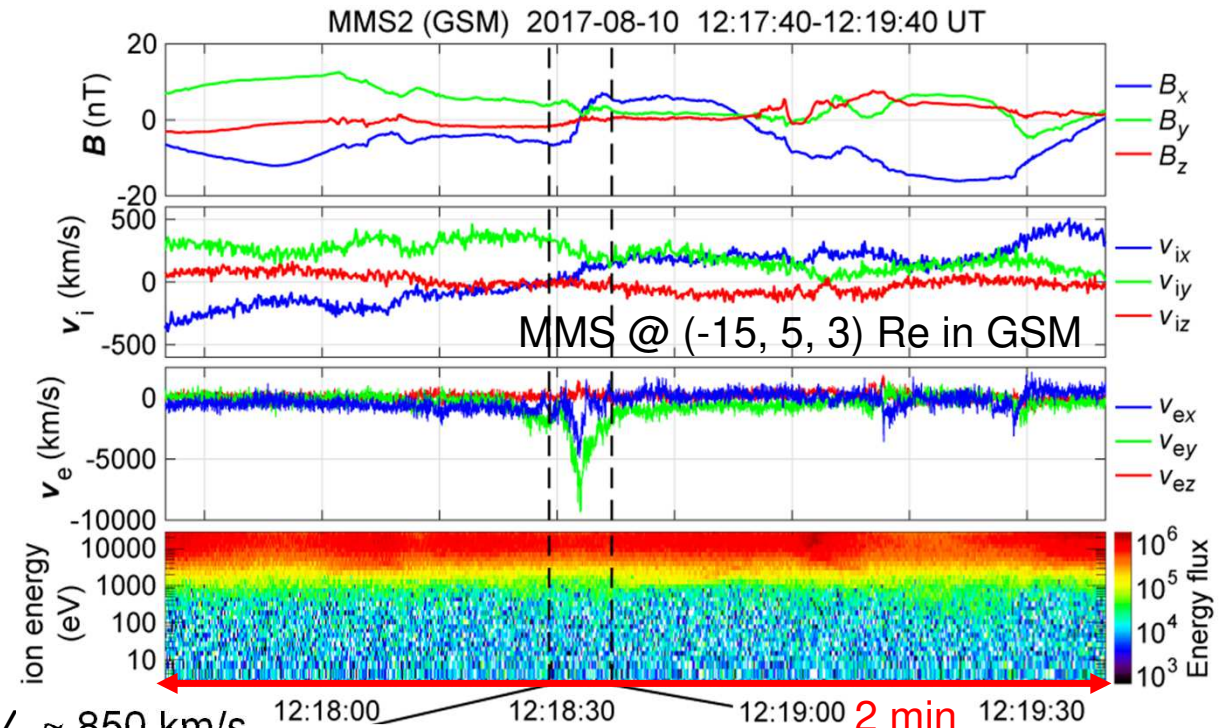
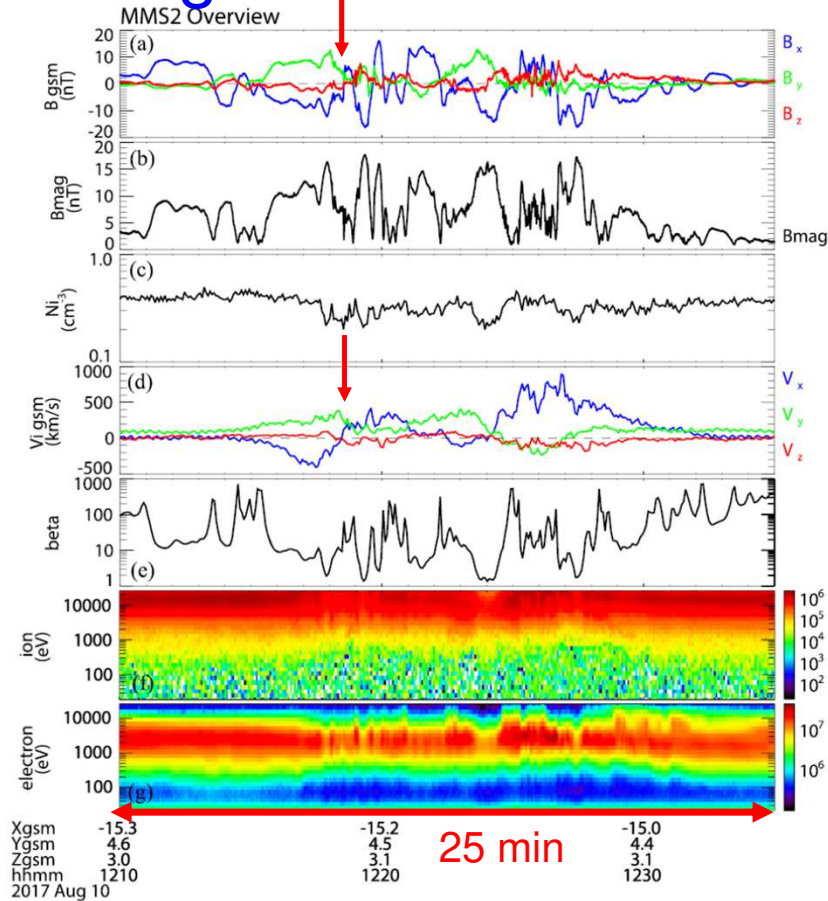
Space Science Institute, Graz

Richard E. Denton

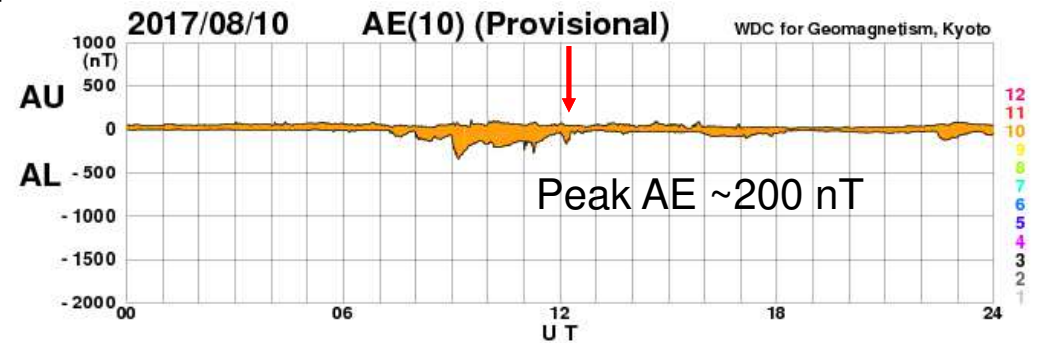
Dartmouth College, Hanover

MMS Fall 2020 Science Working Team Meeting (6-8 October 2020)

# Magnetotail reconnection seen by MMS on 10 August 2017

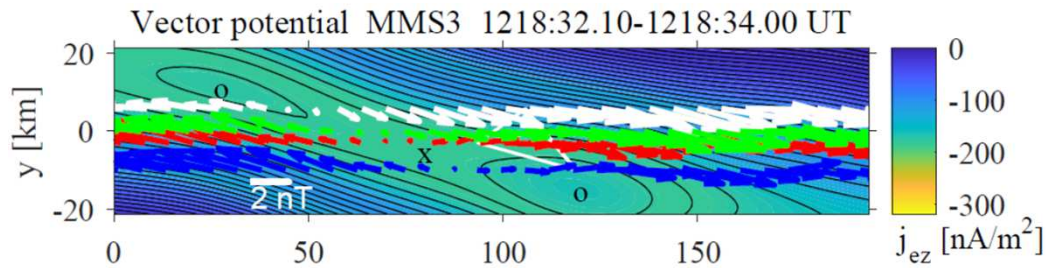


$V_A \sim 850$  km/s  
 $B_0 \sim 15$  nT



- Significant  $\mathbf{j} \times \mathbf{E}$  (Zhou+, ApJ19; Li+, GRL20)
- Crescent-shaped electron distributions
- Super-Alfvénic electron flows

# E-MHD & Polynomial reconstruction results

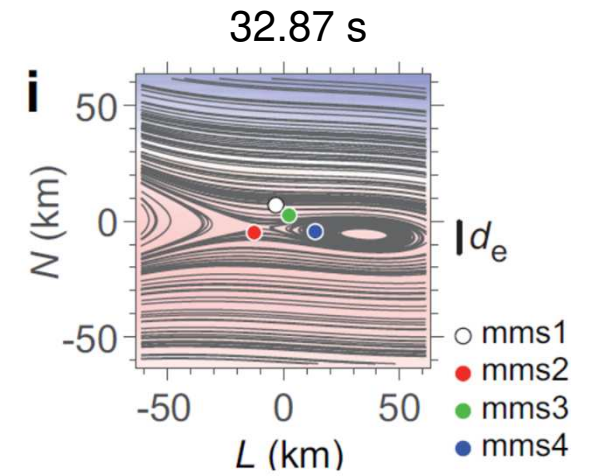
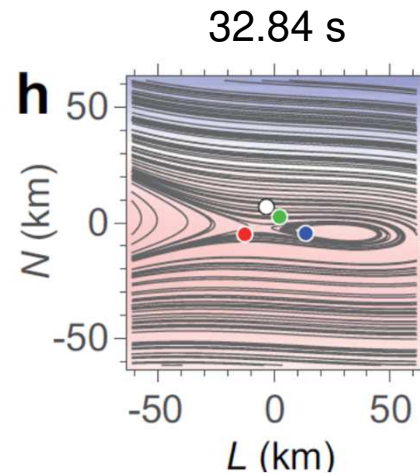
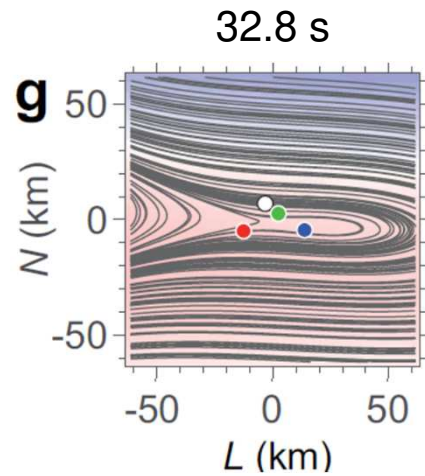
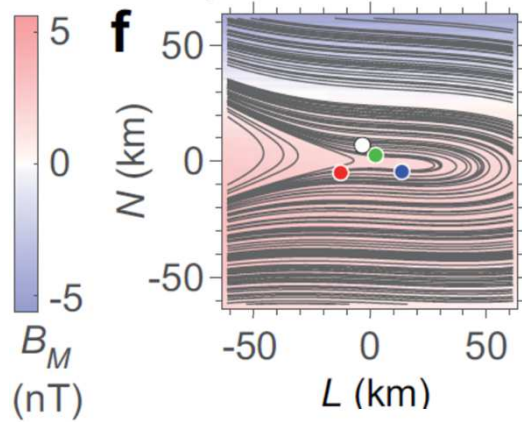
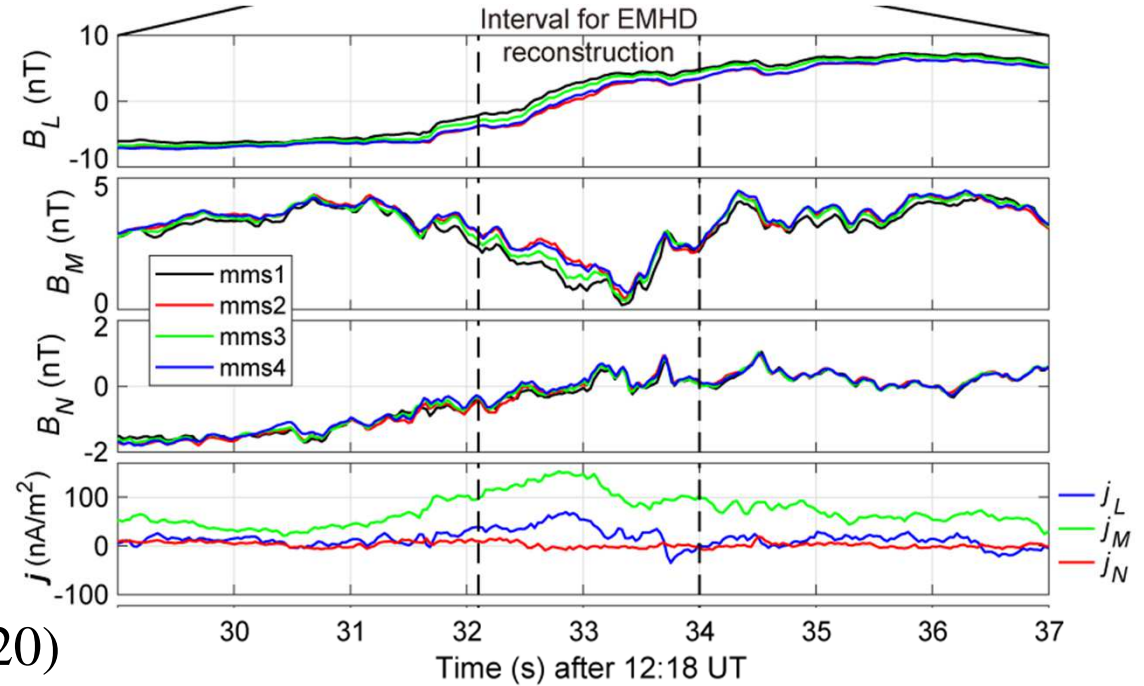


(Sonnerup+, JGR16)

Spacecraft separation  $\sim 18$  km

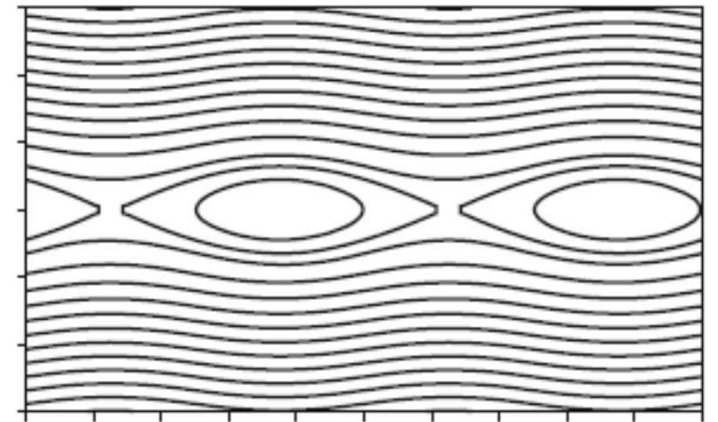
Electron inertial length  $d_e \sim 14$  km (Denton+, JGR20)

Time after 12:18 UT: 32.77 s



## Motivation

- A growing electron-scale magnetic island was identified right in the electron diffusion region (EDR) of magnetotail reconnection (MMS observations of an electron-scale current sheet on 2017-08-10).
- The formation mechanism is not very clear, but may be an electron tearing instability.
- We revisit 2-D tearing instability in electron-scale current sheets, but now including a non-gyrotropic electron pressure effect (Hesse-Kuznetsova dissipation term) (Hesse+, SSR11).





# Linear eigenmode analysis in electron-MHD framework

Generalized Ohm's law 
$$\mathbf{E} = -\mathbf{v} \times \mathbf{B} - \frac{\nabla \cdot \mathbf{P}}{ne} - \frac{m}{e} \left[ \frac{\partial \mathbf{v}}{\partial t} + \nabla \frac{v^2}{2} - \mathbf{v} \times (\nabla \times \mathbf{v}) \right]$$

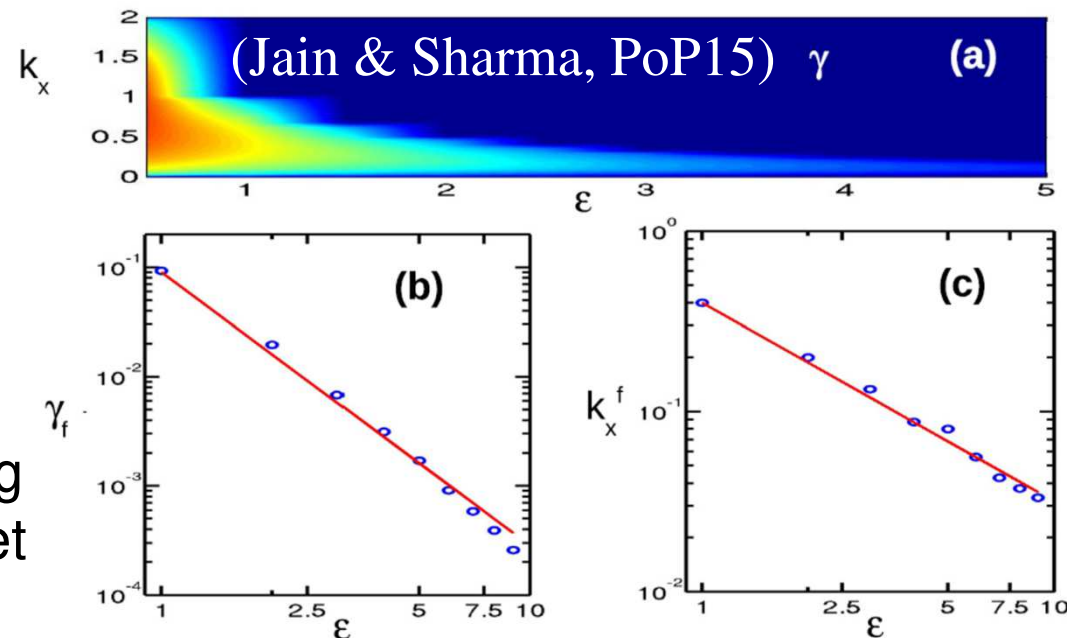
- In the linear eigenmode analysis, Faraday's law is linearized.
- If electrons are isotropic, the electron pressure tensor term does not contribute to the tearing instability.

Faraday's law

$$\frac{\partial}{\partial t} \left( \mathbf{B} - \frac{m}{\mu_0 ne^2} \nabla^2 \mathbf{B} \right) = \nabla \times \left[ \mathbf{v} \times \left( \mathbf{B} - \frac{m}{\mu_0 ne^2} \nabla^2 \mathbf{B} \right) \right]$$

Ampere's law 
$$\mathbf{v} = -\nabla \times \mathbf{B} / (\mu_0 ne)$$

- The growth rate  $\gamma$  (normalized to  $\omega_{ce}$ ) and the wavelength ( $k_x$ ) for the fastest-growing mode are both scaled by the current sheet half-thickness  $\varepsilon$  (normalized to  $d_e$ ).



# Linear eigenmode analysis in electron-MHD framework

- Hesse et al. (SSR, 2011) show that the reconnection electric field in the dissipation region (at the X-point) of steady 2D antiparallel reconnection can be expressed as:

$$E_y = -\frac{1}{n_e e} \left( \frac{\partial P_{exy}}{\partial x} + \frac{\partial P_{eyz}}{\partial z} \right) \approx \frac{1}{e} \frac{\partial v_{ex}}{\partial x} \sqrt{2m_e k_B T_e}$$

- The electron pressure tensor term is then non-negligible in the tearing instability.

Faraday's law 
$$\frac{\partial}{\partial t} \left( \mathbf{B} - \frac{m}{\mu_0 n e^2} \nabla^2 \mathbf{B} \right) = \nabla \times \left[ \mathbf{v} \times \left( \mathbf{B} - \frac{m}{\mu_0 n e^2} \nabla^2 \mathbf{B} \right) \right] - \frac{\sqrt{2mkT_e}}{e} \nabla \left( \frac{\partial v_x}{\partial x} \right) \times \hat{\mathbf{y}}$$

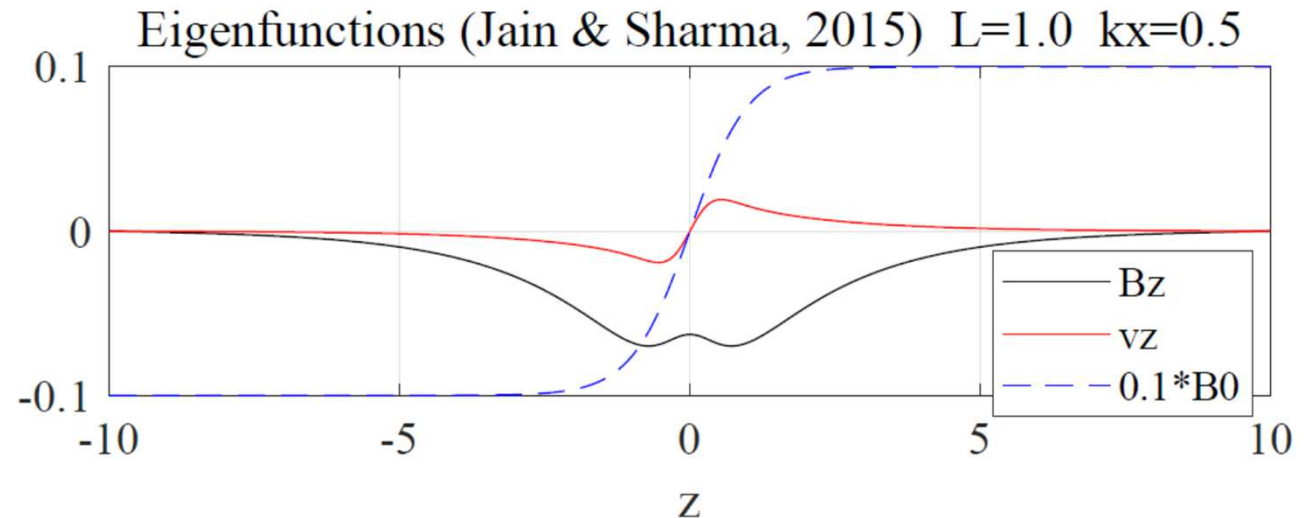
- If normalized to the upstream field intensity  $B_\infty$ ,  $|\omega_{ce}|^{-1}$ , electron inertial length, and electron Alfvén speed,

$$\frac{\partial}{\partial t} (\mathbf{B} - \nabla^2 \mathbf{B}) = \nabla \times [\mathbf{v} \times (\mathbf{B} - \nabla^2 \mathbf{B})] - \sqrt{\beta_e} \nabla \left( \frac{\partial v_x}{\partial x} \right) \times \hat{\mathbf{y}}$$

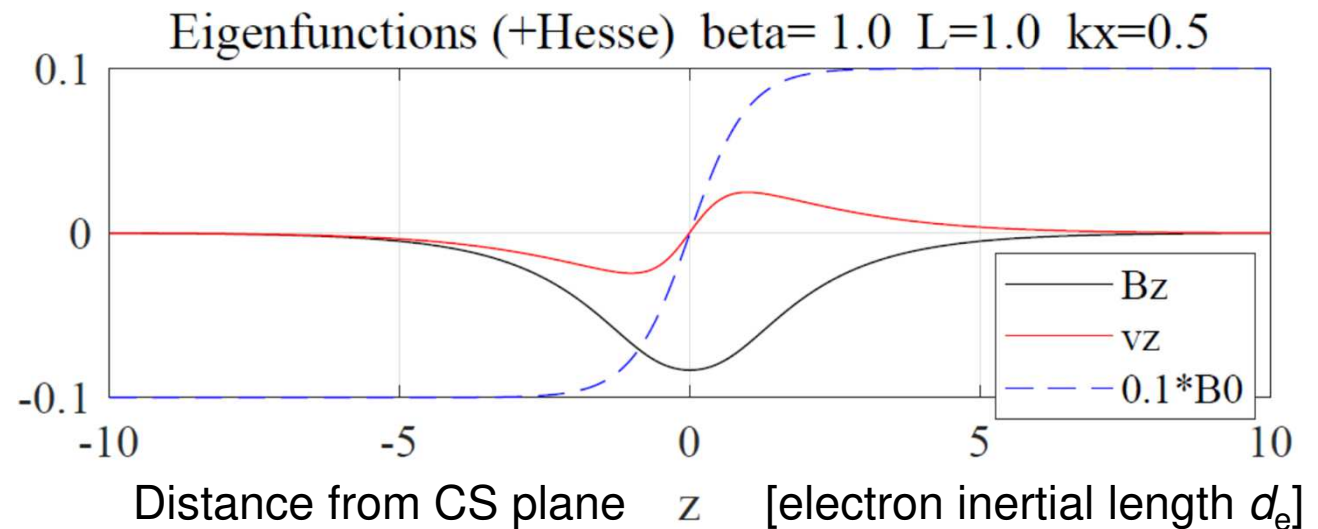
- Assuming uniform density & incompressible perturbation  $\nabla \cdot \mathbf{v} = 0$ , the above equation can be linearized for 2-D perturbations  $\mathbf{B}_1 = \mathbf{B}(z) \exp(ik_x x - i\omega t)$  &  $\mathbf{v}_1 = \mathbf{v}(z) \exp(ik_x x - i\omega t)$  and equilibrium profile  $B_0(z) = \tanh(z/\varepsilon)$ .
- Dependence of the growth rate  $\gamma$  on  $\varepsilon$ ,  $\beta_e$ , and  $k_x$  can be investigated.

# Eigen functions of $B_z$ (reconnected field) & $v_z$ (inflow velocity)

No electron pressure

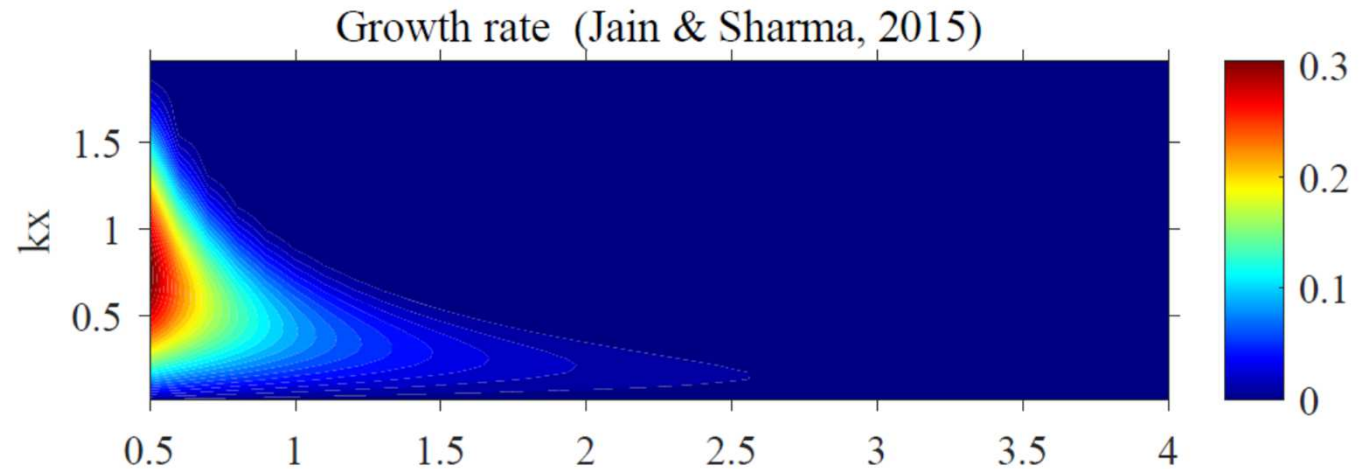


With non-gyrotropic  
electron effect  
(electron beta = 1 case)

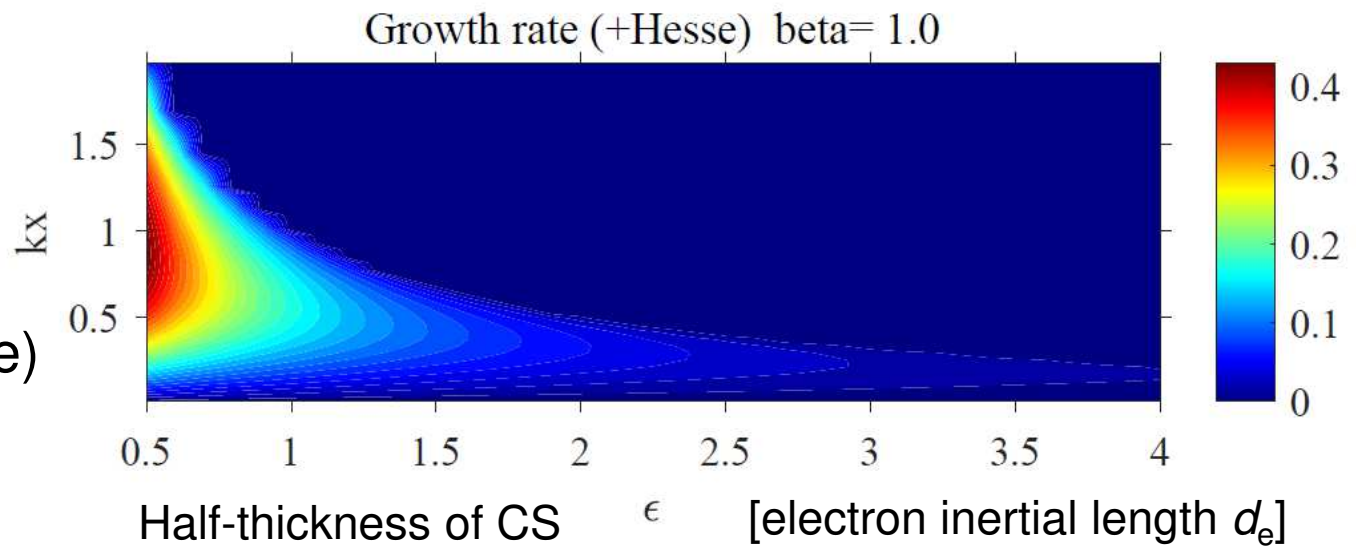


# Growth rate for electron beta = 1 case

No electron pressure



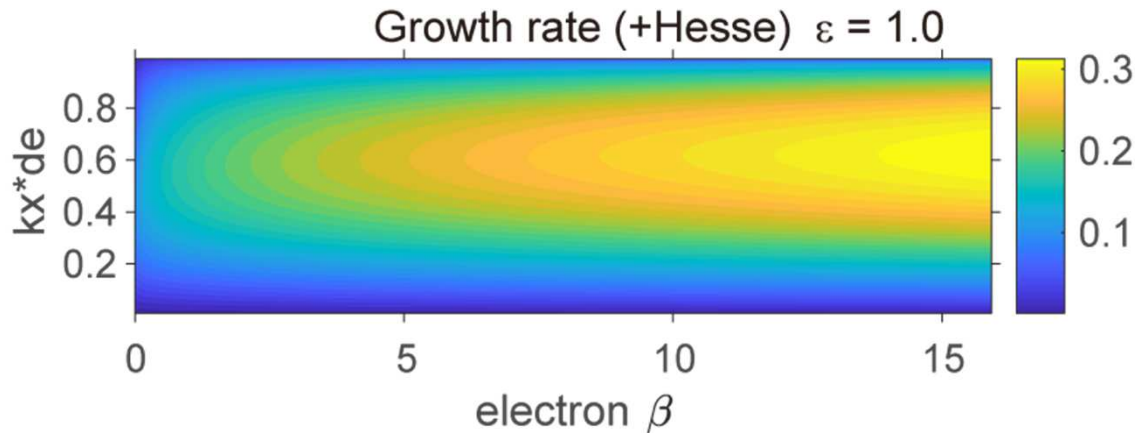
With non-gyrotropic  
electron effect  
(electron beta = 1 case)



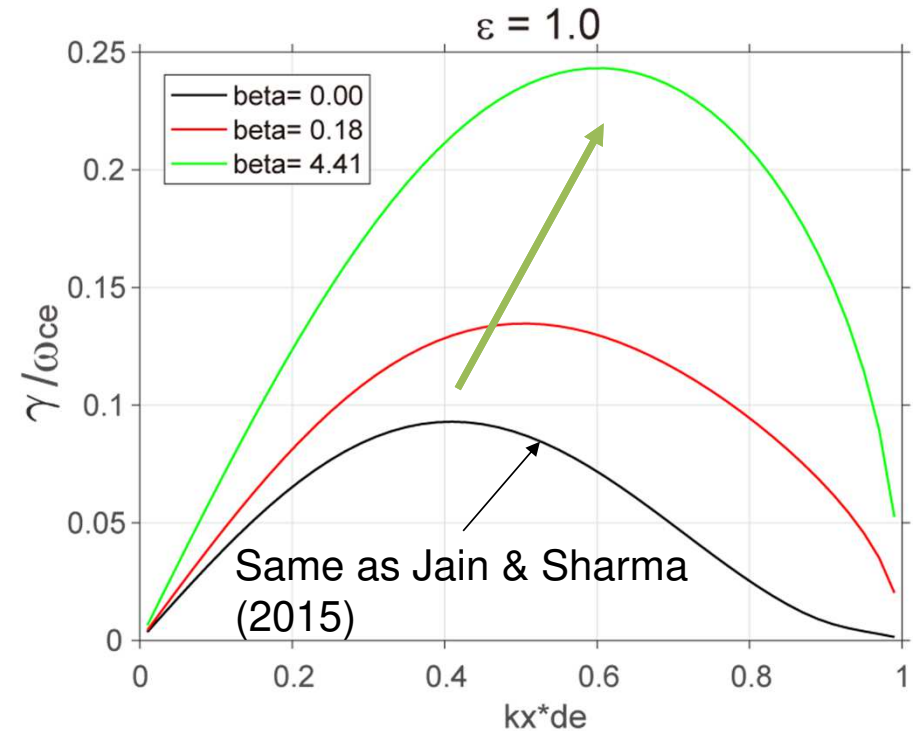


# Electron beta dependence ( $\varepsilon = 1$ case)

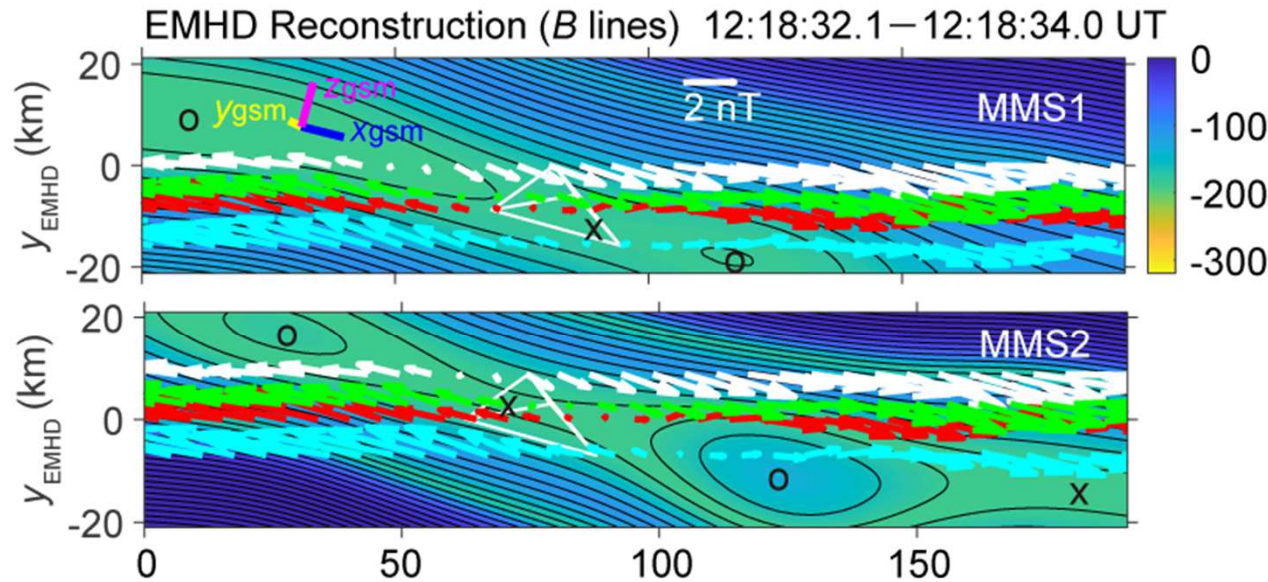
Half-thickness =  $d_e$  case



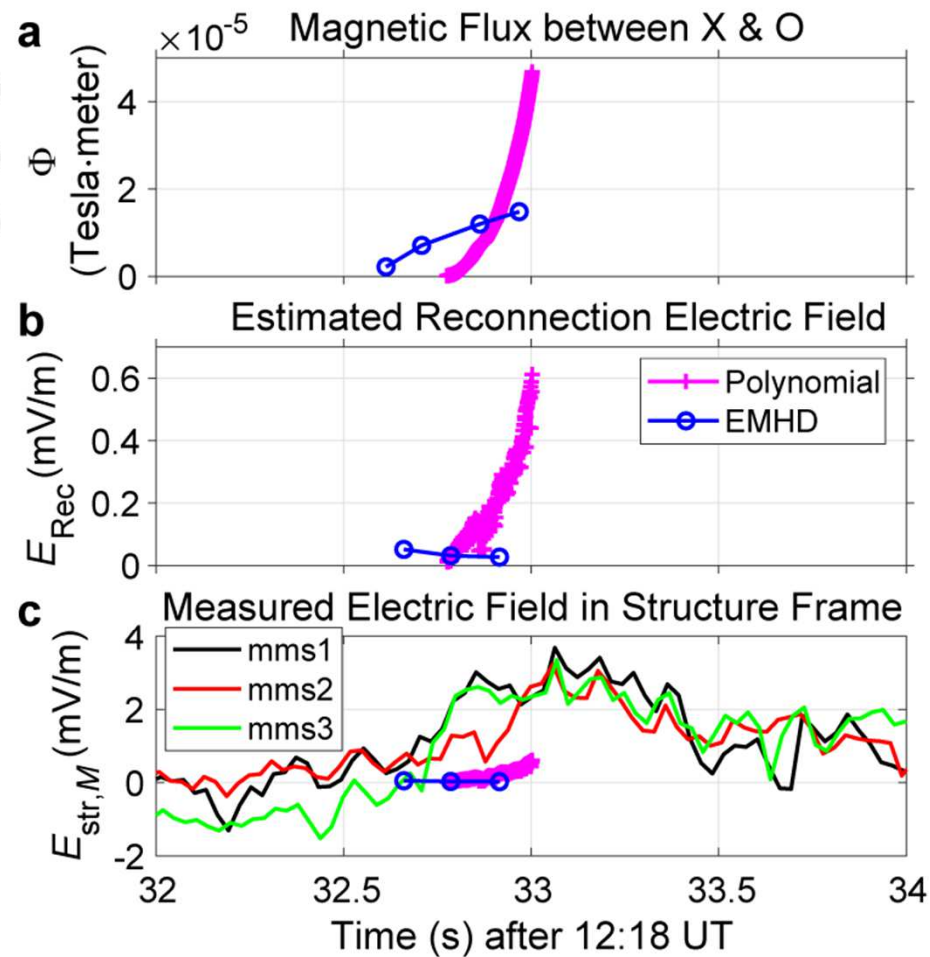
- The growth rate increases with electron  $\beta$ .
- The wavelength for the fastest-growing mode becomes shorter with increasing electron beta.



# Growth time of the observed magnetic island



- In-plane magnetic flux increases with time, consistent with the island growth.
- However, the growth time  $\sim 0.1$  s is far longer than the electron cyclotron time scale ( $2\pi/\omega_{ce} \sim 0.0002$  s), i.e.,
- The growth time is **inconsistent with the tearing instability**.
- The magnetic field in the observed EDR was mostly annihilated, rather than reconnected.



# Magnetic field annihilation in elongated EDR

Ohm's law:

$$\mathbf{E} = -\mathbf{v}_e \times \mathbf{B} + \frac{1}{e} \frac{\partial v_{eL}}{\partial L} \sqrt{2m_e k_B T_e} \mathbf{M}$$

Faraday's law:

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E} = \nabla \times (\mathbf{v}_e \times \mathbf{B}) - \frac{\sqrt{2m_e k_B T_e}}{e} \nabla \left( \frac{\partial v_{eL}}{\partial L} \right) \times \mathbf{M}$$

L-component of Faraday's law:  $B_N \approx 0$

$$\frac{\partial B_L}{\partial t} \approx \frac{\sqrt{2m_e k_B T_e}}{e} \frac{\partial}{\partial N} \left( \frac{\partial v_{eL}}{\partial L} \right)$$

Incompressible flow in/around EDR  $\nabla \cdot \mathbf{v}_e = \frac{\partial v_{eL}}{\partial L} + \frac{\partial v_{eN}}{\partial N} = 0$

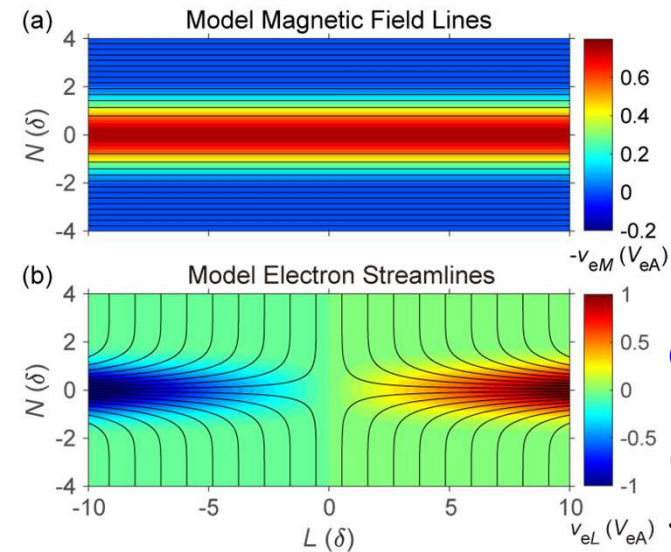
$$\frac{\partial B_L}{\partial t} \approx -\frac{\sqrt{2m_e k_B T_e}}{e} \frac{\partial^2 v_{eN}}{\partial N^2}$$

Using  $-B_\infty v_{eN} \approx B_L V_\infty$  (right figure), we derive a diffusion equation for  $B_L$

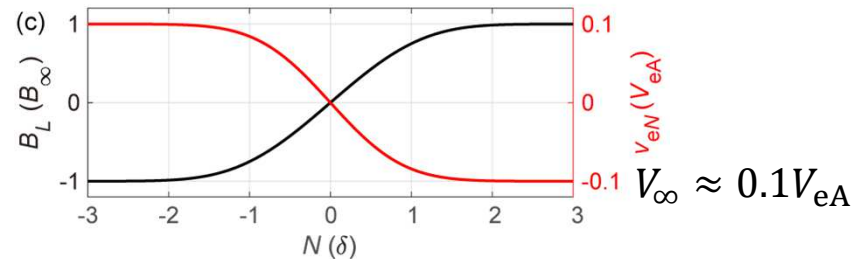
$$\frac{\partial B_L}{\partial t} \approx D_B \frac{\partial^2 B_L}{\partial N^2}$$

$$\text{Diffusion coefficient: } D_B = r_{ge} V_\infty$$

- In an elongated EDR, the electron pressure tensor term can act for magnetic field annihilation.
- The diffusion coefficient  $D_B$  is proportional to the inflow speed  $V_\infty$  and electron gyroradius  $r_{ge}$ .



Exact solution of dissipative EMHD equations (Sonnerup+, JGR16)



## Summary & Discussion

- The linear eigenmode analysis of the 2-D tearing instability with non-gyrotropic electron pressure effect (Hesse-Kuznetsova term) suggests
  - The growth rate increases with increasing electron beta  $\beta_e$ .
  - The wavelength for the fastest-growing mode is shorter for larger  $\beta_e$ .
- However, the rate of magnetic-field annihilation in an extended electron-scale current sheet may also be enhanced with increasing  $\beta_e$ .
- The tearing instability and magnetic-field annihilation might be competitive in electron-scale current sheets.
- Is the EDR width scaled by the electron inertial length or gyroradius?
  - If scaled by gyroradius ( $\beta_e$ ), the tearing would be faster for lower  $\beta_e$  while the annihilation may not depend much on  $\beta_e$  (larger EDR width for larger  $\beta_e$ ).
  - Any  $\beta_e$  dependence of either or both processes in simulation & observations?