

The background features a stylized sun at the bottom center, with a color gradient from yellow to red. Numerous yellow lines with arrowheads radiate from the sun, representing magnetic field lines and particle beams. The sky is a dark blue-purple gradient. At the bottom, a greenish-yellow ground line is visible with several blue circles representing particles or ions.

# Recent Advances in Electron Beam Transport and Implications for White-Light Flares

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**Sun-Climate Symposium**

# Overview of Talk

White-light (optical and NUV continuum) radiation

Electron beams in solar and stellar flares

**The problem of high energy fluxes**

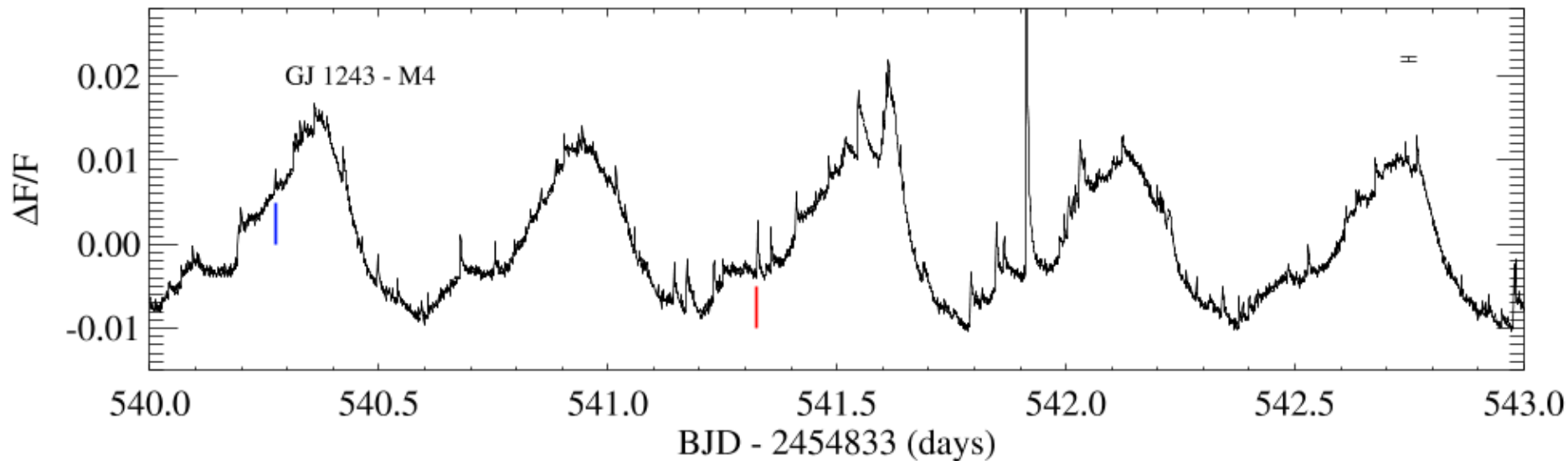
**Overview of Kontar et al. 2012 beam-plasma transport theory and results**

**Application to white-light & line emission in stellar flares (Kowalski 2023  
*ApJ Letters*)**

New directions: some first results

# White-Light Stellar Flares

Broadband optical continuum enhancements observed on cool stars with Kepler, K2, TESS



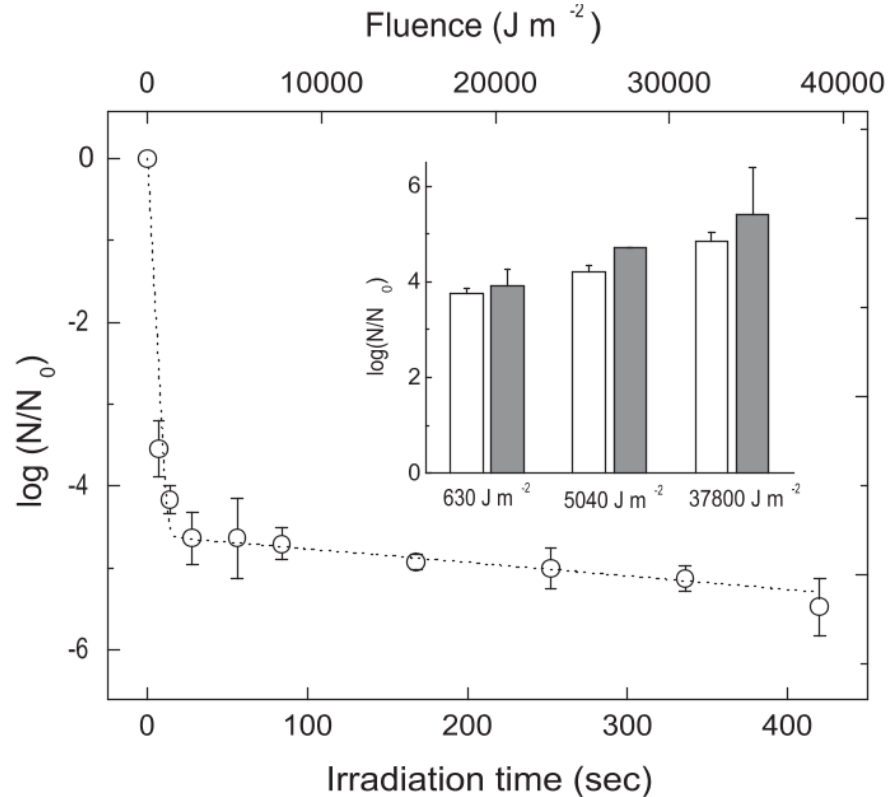
Hawley et al. 2014 (Kepler data of GJ 1243 -- active M4 main-sequence star)

11 mo of 1-minute cadence data

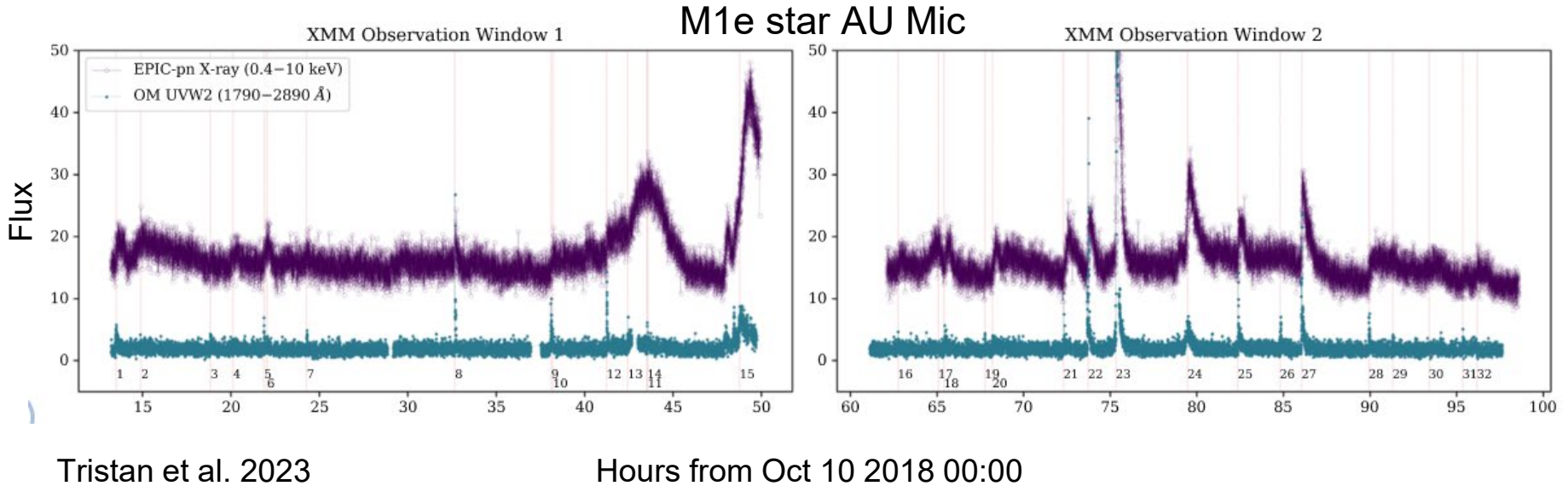
# Biological Effects in M-dwarf Habitable Zones

Laboratory experiments of 2000-2800 Å superflare irradiation of an exoplanet (Abrevaya et al. 2020; right)

X-rays from flare can be reprocessed into NUV (Smith et al. 2004)



# Flares are faster in NUV, slower in SXR

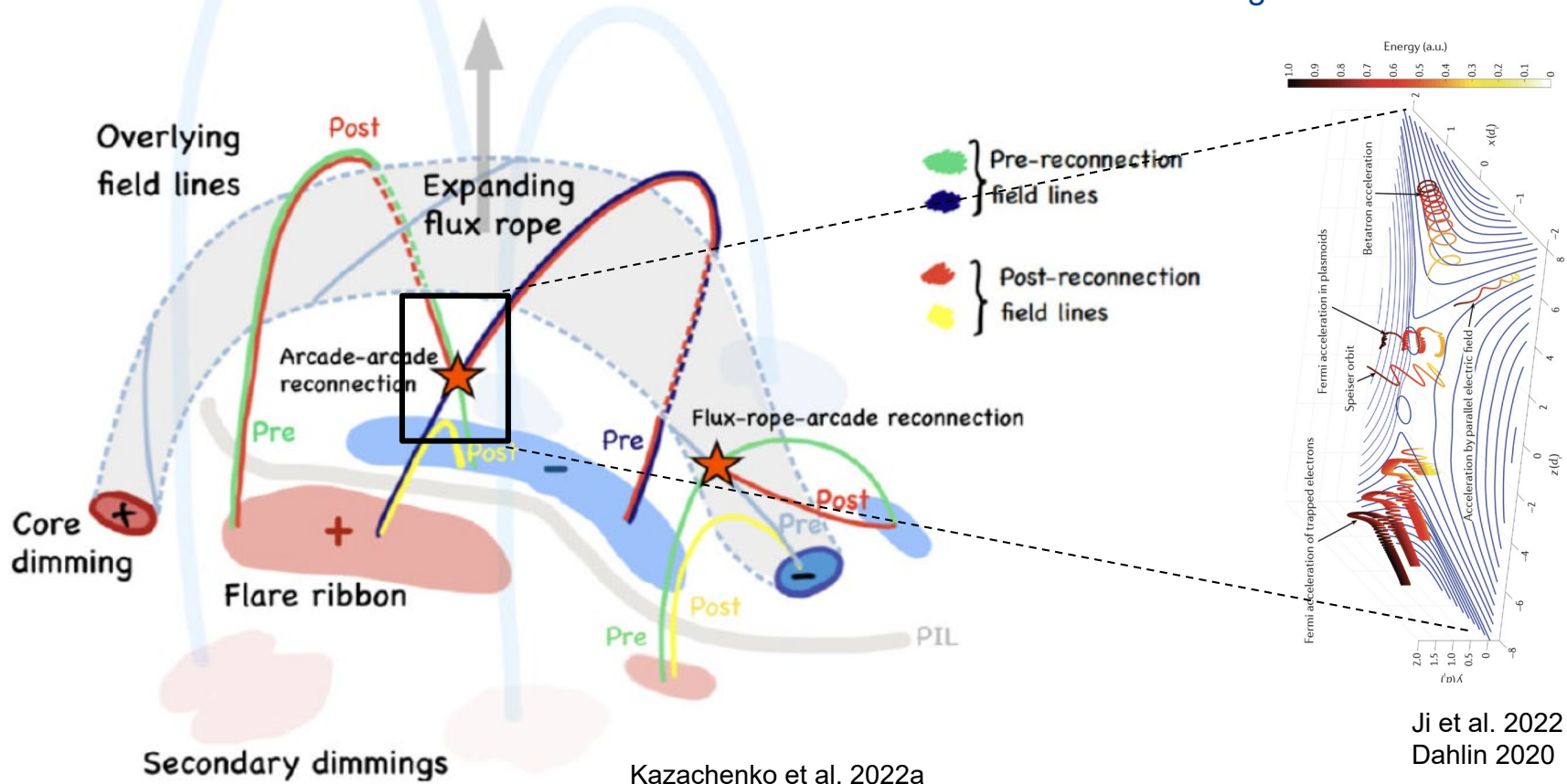


$$2/3 E_{\text{SXR}} \sim E_{\text{U-band}} \sim 1/10 E_{\text{Bol}}$$

See also Osten & Wolk 2015

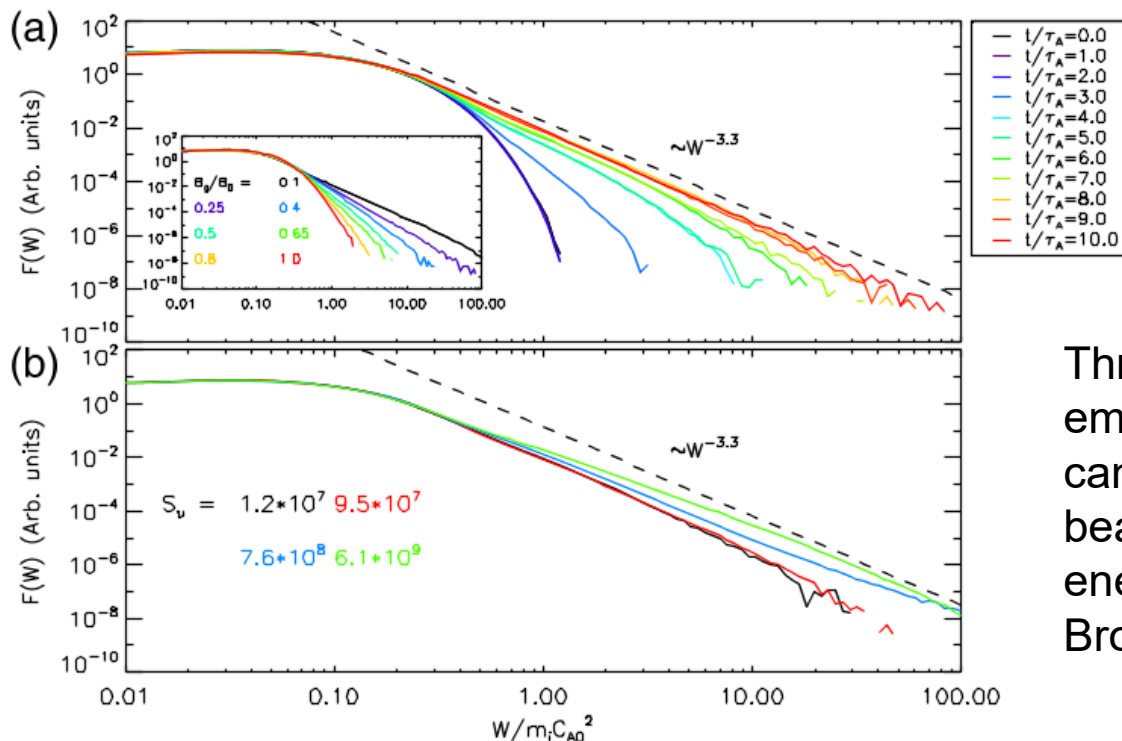
# Reconnection **under** erupting filament and **with** erupting filament

Particle acceleration & transport are key in **linking** magnetic field action to heating/radiation



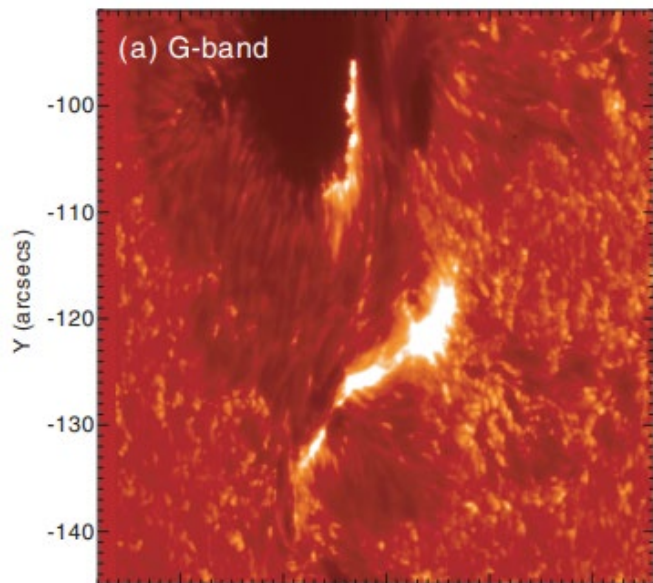
# Power-laws of electron beams produced in corona

~ 5-20% of background plasma accelerated into power-laws

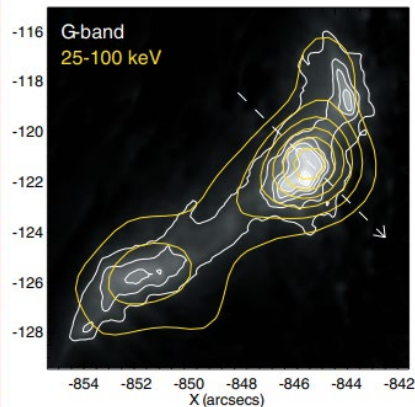


Through hard X-ray brems emission in chromosphere, we can infer the properties of the beams assuming collisional energy loss only (“CTTM” theory; Brown 1971)

# Large Nonthermal Electron Fluxes ( $\text{erg/s/cm}^2$ ) into the Chromosphere



Krucker et al. 2011



Standard collisional thick target modeling (CTTM) infers fluxes of  $10^{12} - 10^{13} \text{ erg/s/cm}^2$  above  $E_{\text{cutoff}} = 18\text{-}20 \text{ keV}$

*25-100 keV: Non-thermal Brems hard-X-rays (from RHESSI)*

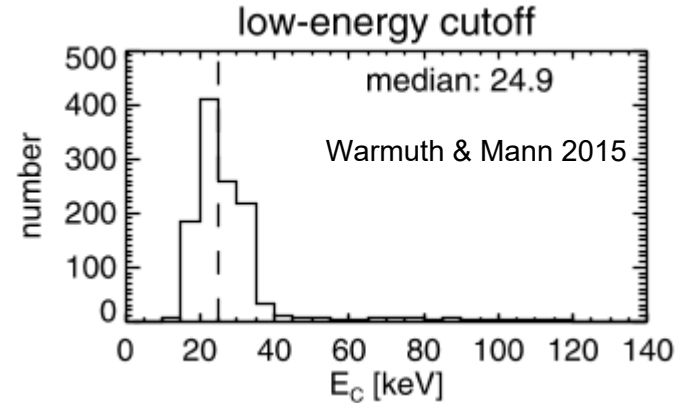
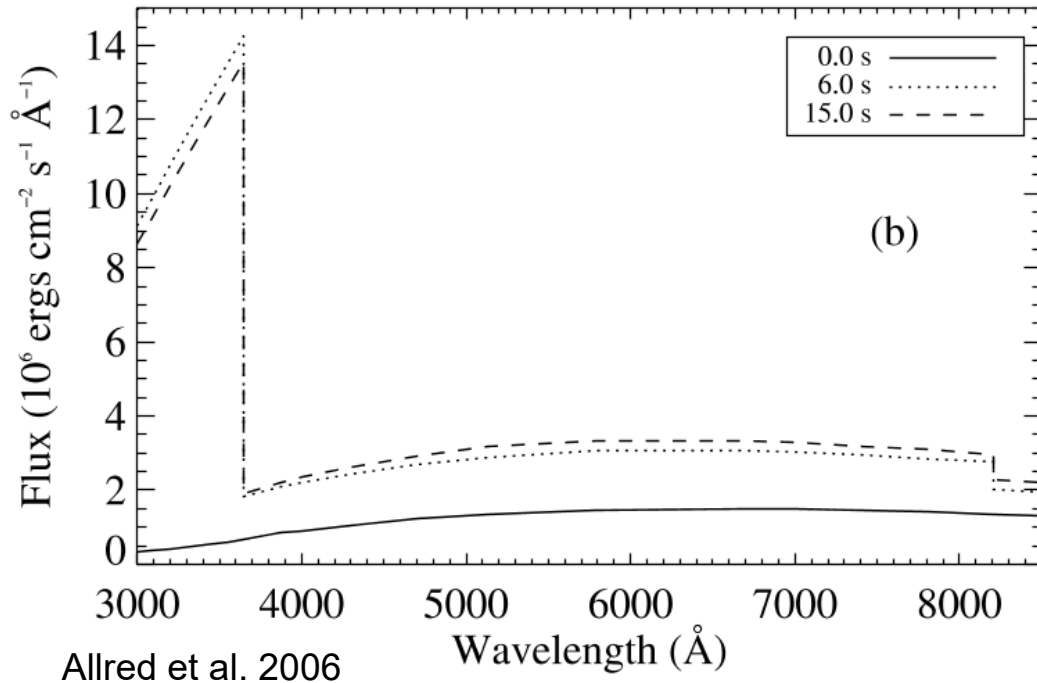
Another flare (Mar 29 2014):  $3.5 \times 10^{11}$  (“3.5F11”; Kleint et al. 2016) to  $\sim 2 \times 10^{12} \text{ erg/s/cm}^2$  (“2F12”; Kowalski et al. 2017) above  $\sim 20 \text{ keV}$  inferred.

**Long story short:** beam particles should **thermalize** in corona through return current electric fields and beam-plasma instabilities... Only catastrophic energy loss = **DOOM** of coronal electron beams ?!



# Solar-type Electron Beams with small low-energy cutoffs

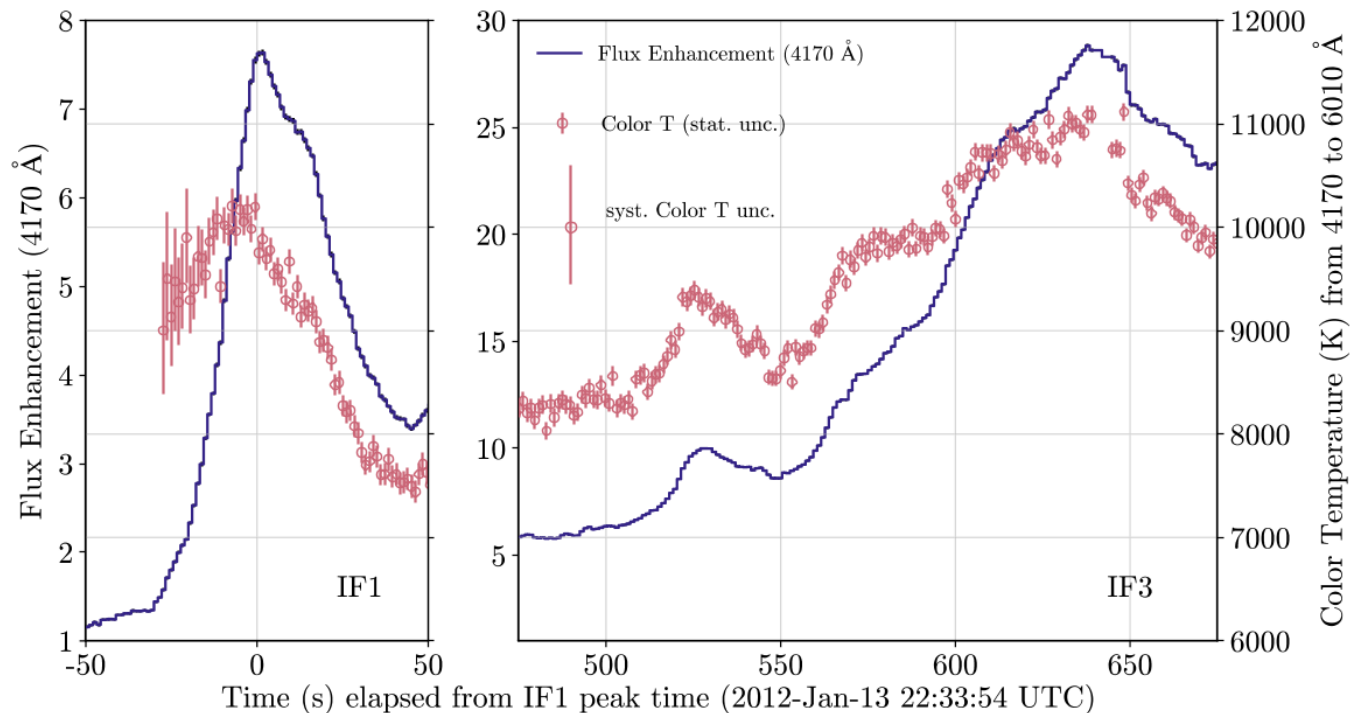
Models of M-dwarf flare heating using electron beam parameters derived from “CTTM” (collisional thick target model) of hard X-ray footpoints (Brown 1971)



$$f(E) = f_o \left( \frac{E}{E_{\text{cutoff}}} \right)^{-\delta}$$

# Giant Flare from YZ CMi

## Observational evidence for deep heating in stellar (M dwarf) flares



Nothing new to see here, folks, except **high-time res of T ~ 9000 - 11,000 K** color temp in rise and peak phases (Hawley & Fisher 1992, Fuhrmeister+08, Kowalski+2013, ...)

ULTRACAM data from Kowalski+16, filter ratios (**4170/6010Å**) converted to color temps in K2023

What physical processes could (possibly) explain extreme beams -- **and the deep heating** -- in stellar flares?

Thankfully, I didn't have to ask ChatGPT.

I thank Eduard Kontar for pointing me in a productive direction and for helpful discussions.

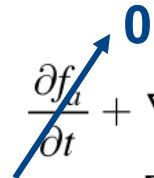
# Modeling nonthermal particle transport in an atmosphere

## A: Particle-in-cell (e.g., Li, Drake, & Swisdak 2012)

- collisionless (Vlasov-Maxwell), very short temporal & spatial scales

## B: Beam distribution function in phase space (Allred+2020)

- collisional Fokker-Planck treatment, time-indepen
- able to calculate background plasma heating


$$\frac{\partial f_a}{\partial t} + \nabla_r \cdot (\mathbf{v}f_a) + \nabla_p \cdot (\mathbf{F}^e f_a) + \sum_b \nabla_p \cdot \mathbf{j}^{a/b} = 0,$$

## C: Time-dependence of background plasma wave energy density (Kontar+2012)

- the “quasilinear” / weak turbulence theory
- Coulomb collisions included but other simplifications made

## Option C: Time-Dependence and Background Plasma Evolution (Langmuir and Ion-Acoustic Waves):

Kontar et al. 2012 solves time-dependent distribution function ( $t = 0$  to 1s) with background plasma waves and with Coulomb collisions

$$f(v, t = 0) = \frac{n}{\sqrt{2\pi}v_{Te}} \exp\left(-\frac{v^2}{2v_{Te}^2}\right) + g_0(v),$$

$$1) \quad \frac{\partial f}{\partial t} = \frac{4\pi^2 e^2}{m^2} \frac{\partial}{\partial v} \left( \frac{W_k}{v} \frac{\partial f}{\partial v} \right) + St_{col}(f),$$

$$2) \quad \frac{\partial W_k}{\partial t} - \frac{\partial \omega_{pe}(x)}{\partial x} \frac{\partial W_k}{\partial k} = \frac{\pi \omega_{pe}^3}{nk^2} W_k \frac{\partial f}{\partial v} - \gamma_{col} W_k + \frac{\omega_{pe}^3 m_e}{4\pi n_e} v \ln\left(\frac{v}{v_{Te}}\right) f + St_{decay}(W_k, W_k^s).$$

Low-energy electrons lost fastest due to Coulomb collisions (bump-on-tail)  $\Rightarrow$  background electron (Langmuir) waves

$W_k$  : Langmuir waves

$W_k^s$  : ion-sound waves (ion-acoustic)

**Critically, includes 3-wave langmuir<sub>k1</sub>  $\Leftrightarrow$  langmuir<sub>k2</sub> + ion-sound**

## Summary of Kontar et al. 2012 results

Solved simultaneous equations of (1) distribution functions for Maxwellian backgd + Powerlaw beam (for  $n_{\text{beam}} / n_{\text{backgd}} \sim 1e-2$ ) and (2) backgd plasma waves over  $\Delta t = 1$  s

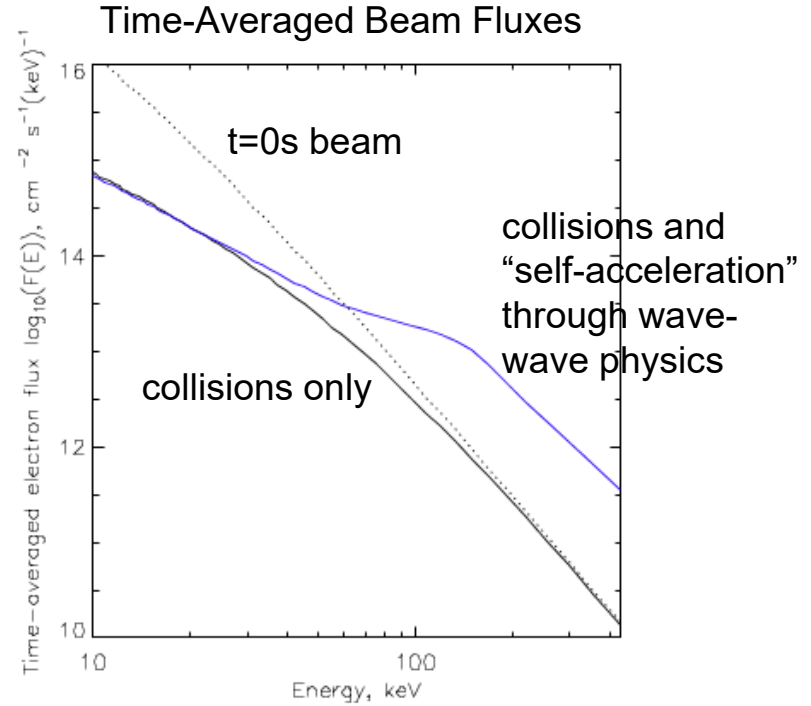
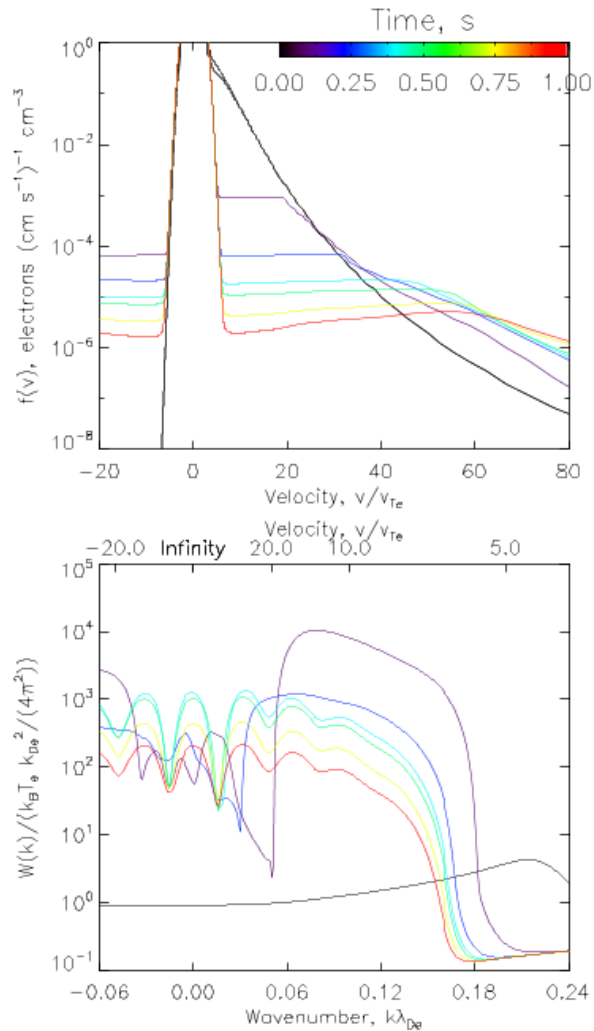
Collisional loss terms (drag and diffusion) on beam electrons (simplified Fokker-Planck)

Regimes not accessible in PIC simulations (but have been investigated with PIC)

Ion-acoustic and Langmuir plasma waves: diffusion, refraction, and **non-linear wave-wave interactions (decay and scattering)**

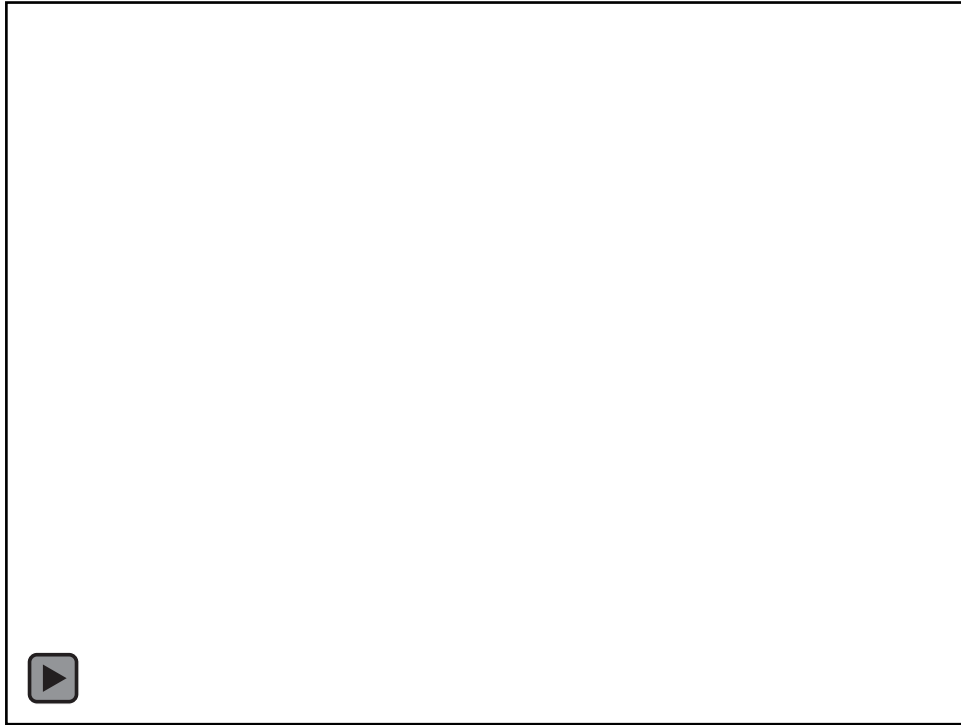
Background density fluctuations and density gradient (important too)

# Kontar et al. 2012 A&A: Main Results



**Fig. 5.** Same as Fig. 2 but with ion-sound wave interactions as well as density fluctuations.

# Langmuir Wave Demo (mono-k)



$$v_{\text{phase}} = \frac{\omega_{pe}}{k(t)} \approx 0.5c$$

k changes smoothly through diffusion & refraction

k decreases abruptly through non-linear 3-wave processes

resonant acceleration of beam electrons (here, 80 keV) with phase velocity of electric fields in wave (~Landau damping)

$$T = 10^6 \text{ K} \quad n_e = 5 \times 10^{10} \text{ cm}^{-3} \quad \omega_{pe} \approx 10^{10} \text{ rad s}^{-1}$$



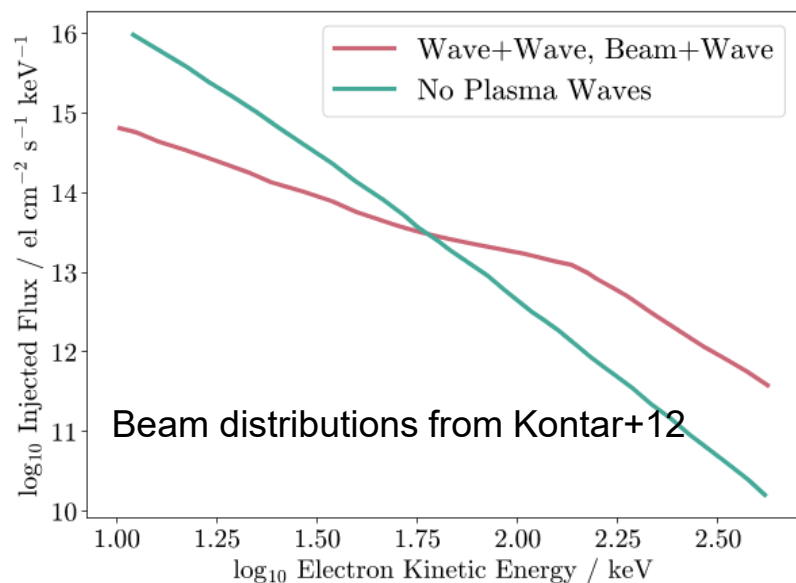
## Story so far

**Kontar et al. 2012:** postulated that increase in  $E > 100$  keV electrons systematically leads to higher production of hard X-rays with far fewer number of total electrons; possibly alleviates problems from large fluxes inferred using CTTM, as in Krucker et al. 2011

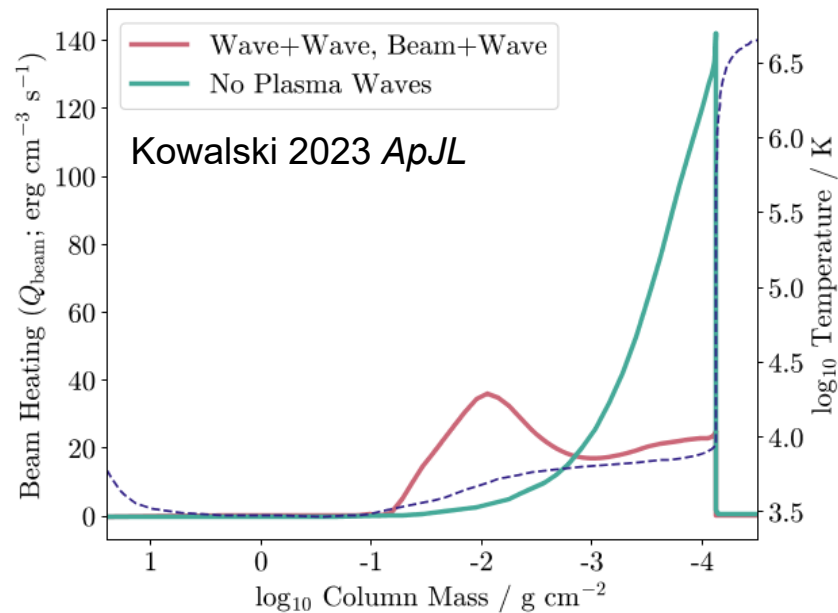
**For stellar flares (Kowalski 2023):** With a 10x larger number of  $E > 100$  keV electrons than typically assumed from CTTM power-laws in solar flares, what are the implications for heating the low chromosphere and producing NUV and optical continuum radiation?

# Simulated M-dwarf chromospheric heating to Kontar et al. 2012 beam

Use RADYN and Fokker-Planck solution (Allred et al. 2015)

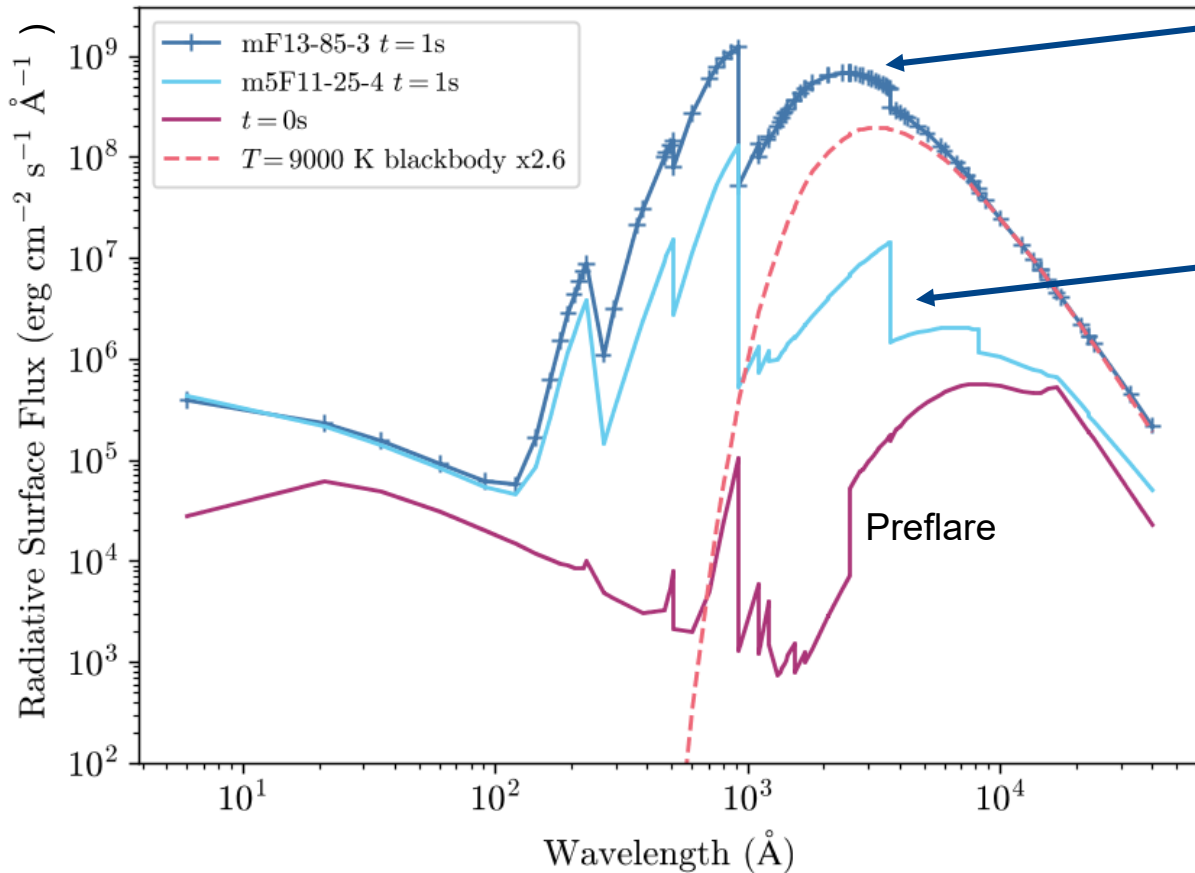


(a)



(b)

# Heating in low vs. high chromosphere



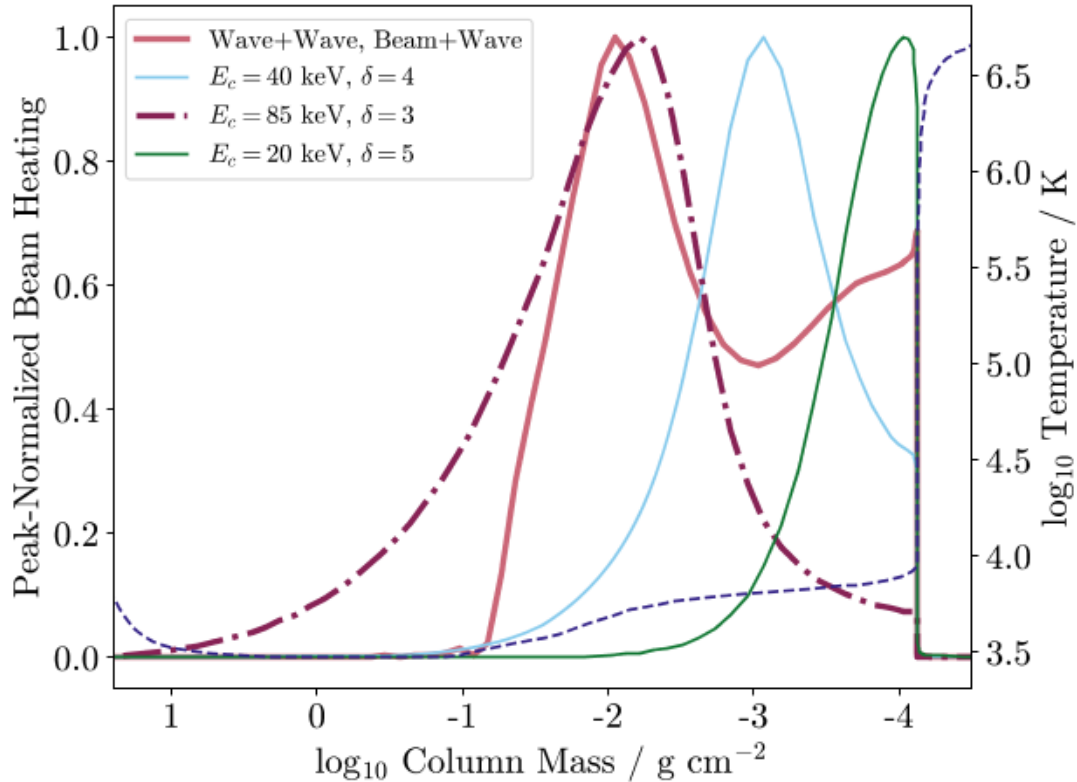
Heating  $m_{\text{col}} \sim 0.01$   
 $\text{g/cm}^2$  to  $1e4$  K

Heating  $m_{\text{col}} \sim 0.001$   
 $\text{g/cm}^2$  to  $1e4$  K

Preflare

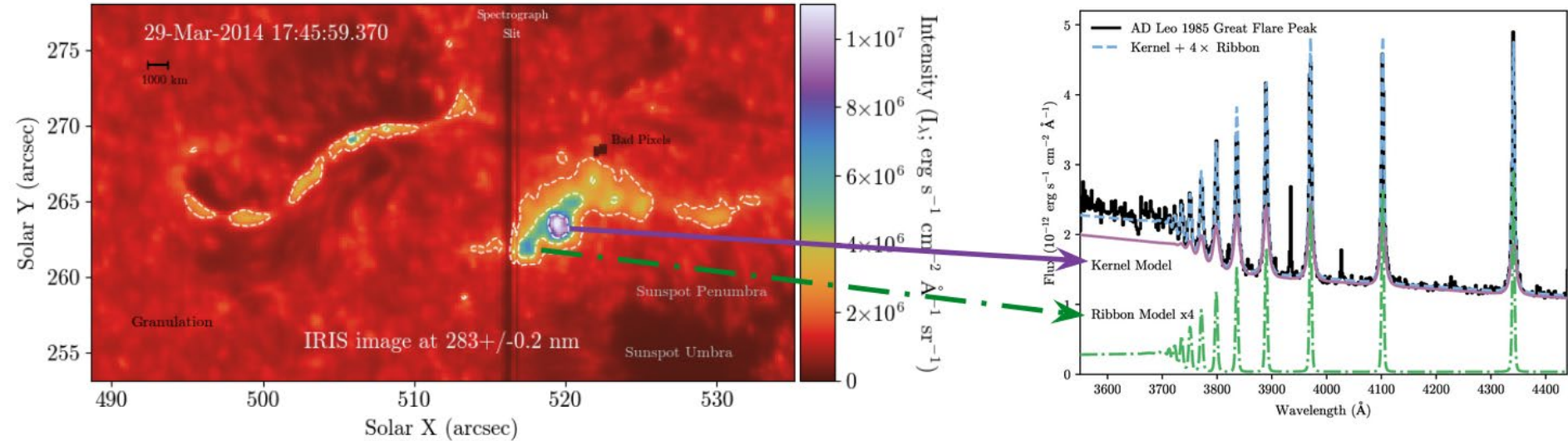
Calculations from RADYN  
(Carlsson & Stein 1995/97)

# Compared to Heating Rates from Large $E_{\text{cutoff}}$



Heating rate maximum around  $0.01 \text{ g/cm}^2$  column mass similar to *semi-empirical forward modeling* approach with large,  $E_{\text{cutoff}} = 85 \text{ keV}$  and  $\delta = 3$  RADYN models (Kowalski+17, 2022) and in upcoming public grid

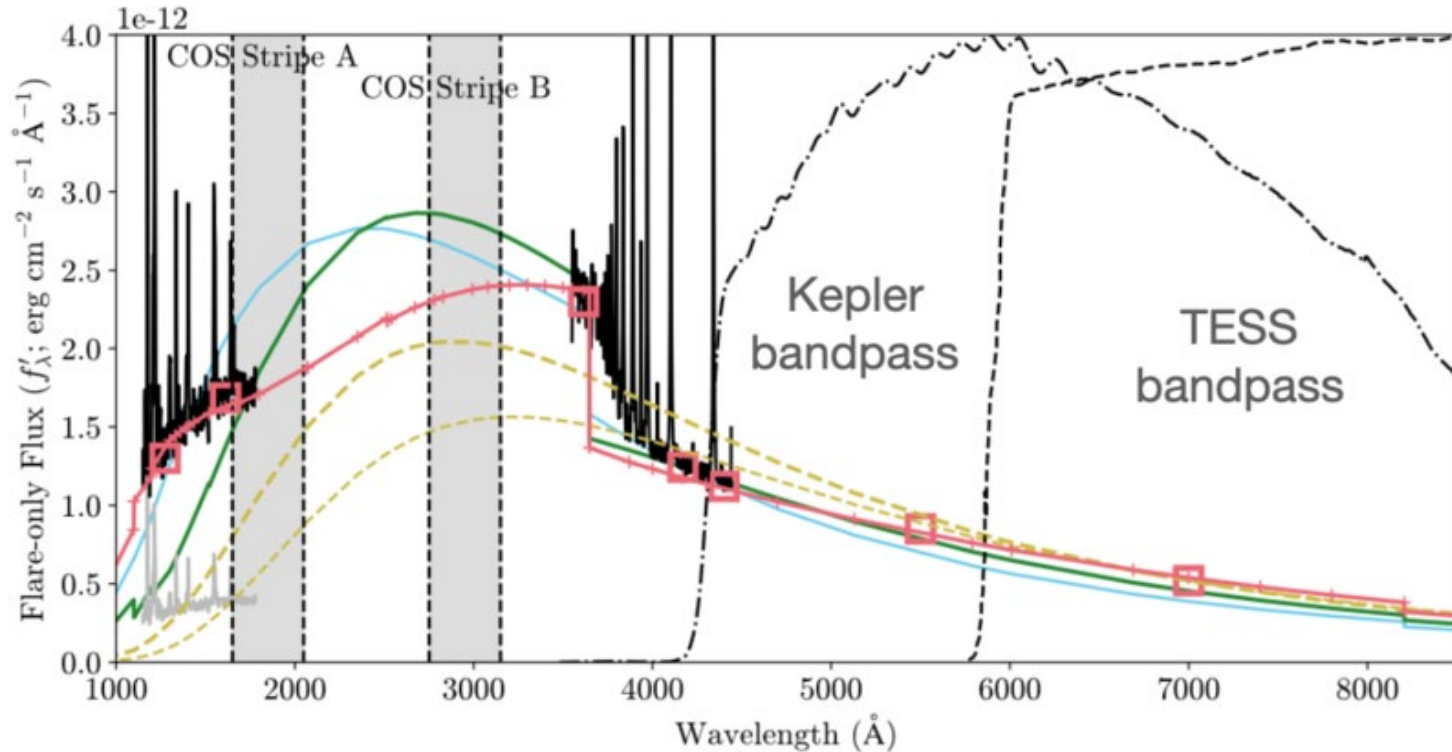
# Kernels: most of wing broadening and optically thick $T \sim 1e4$ K continuum (with Large $E_{\text{cutoff}}$ Model)



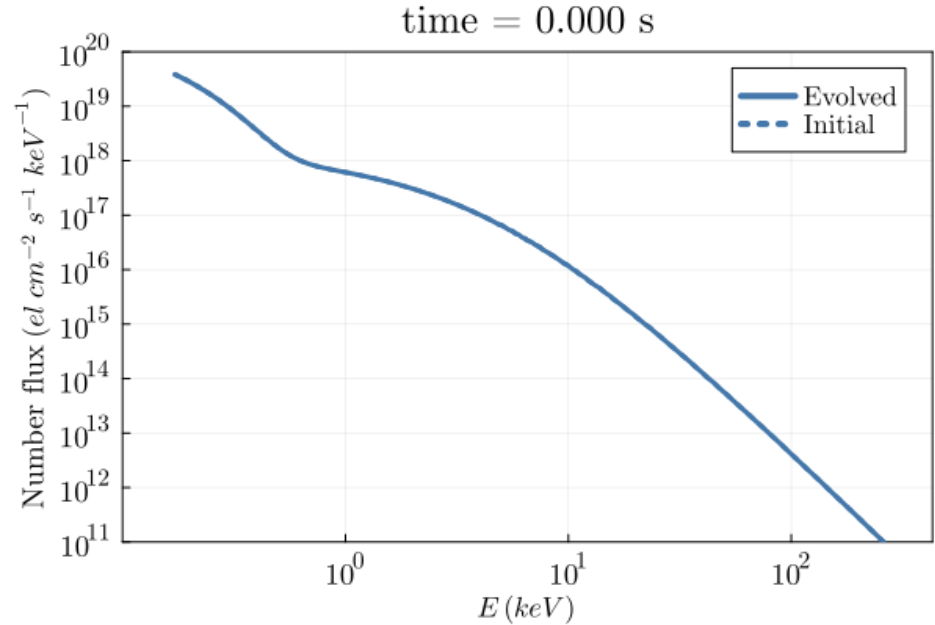
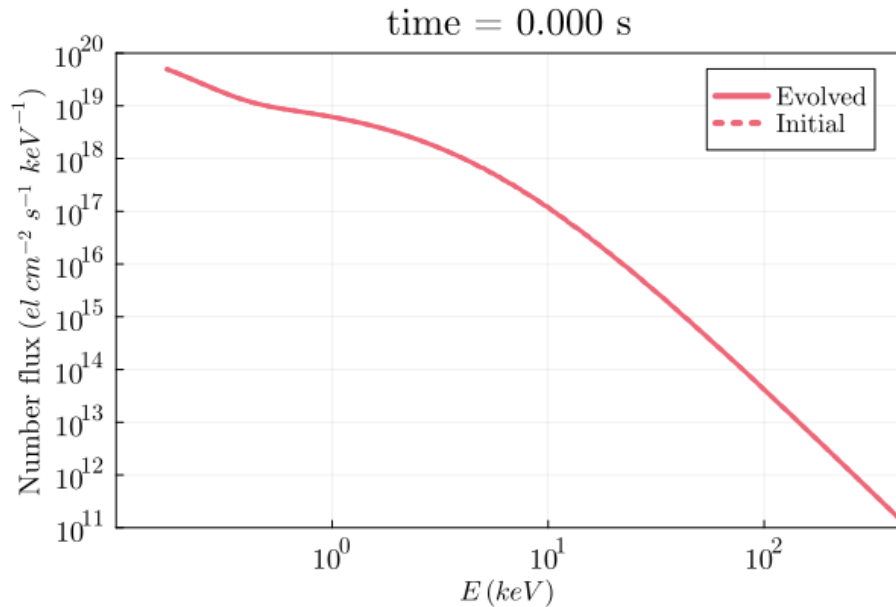
IRIS SJI 2832 image of 2014-Mar-29 solar flare

A superflare ( $10^{34-5}$  erg) event on AD Leo from Hawley & Pettersen 1991

# Can use good models to predict missing wavelengths



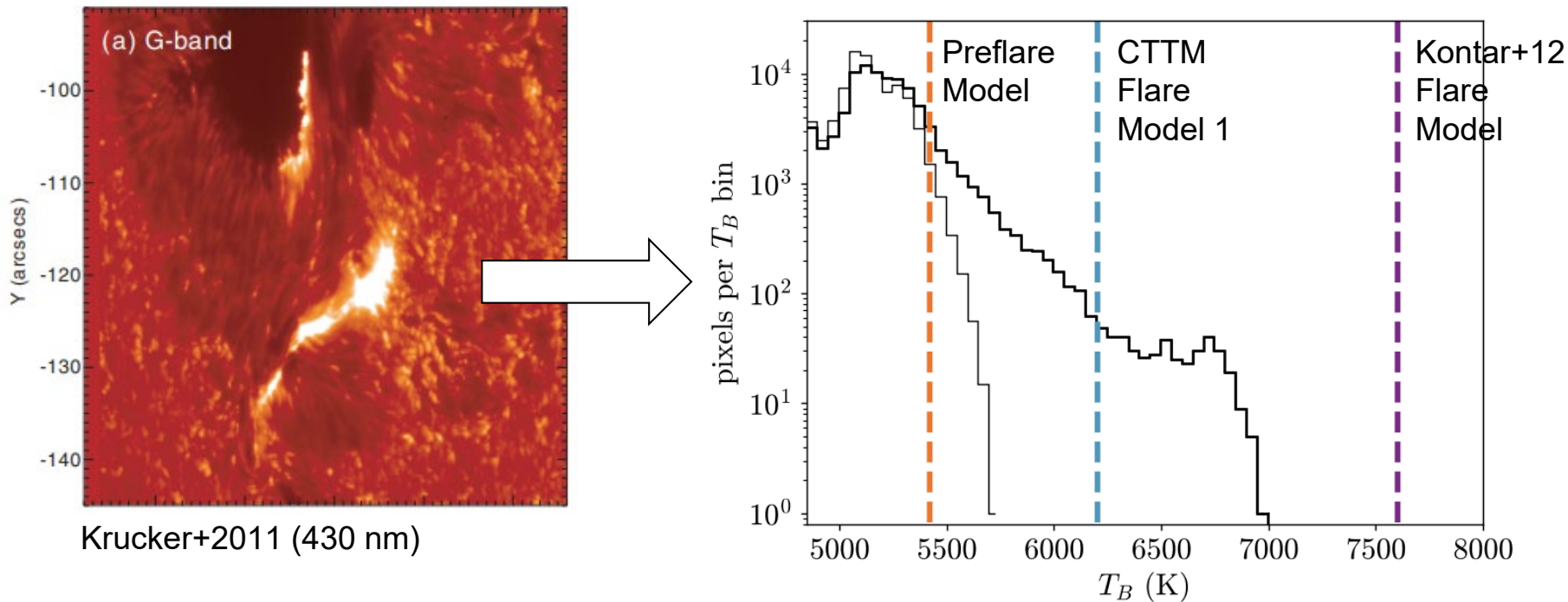
# Much parameter space to explore with Kontar+12 code!



New simulation (left) with larger beam density

*Note: still some problems to work out (short-timescale effects, fully relativistic theory d.n.e.)*

# New interpretations of solar flare optical and hard X-ray data



New radiative-hydro RADYN calculations driven by Kontar+12 beam heating



# Conclusions

The **beam transport and plasma wave interaction** theory of Kontar et al. 2012 produces an **enhancement in  $E > 100$  keV electrons** that may plausibly provide a physical explanation for the **deep heating in stellar flares around  $\log \text{col mass} \sim -2$**

*An alternative hypothesis* to CTTM-inferred power-laws is now possible using RADYN simulations of solar and stellar flares

Many other problems in solar / stellar flares can be investigated (interpretation of HXR data, multi-wavelength energy budgets, anomalies in radio spectra, etc...)